

"Integrability in Gauge and String Theory, 22 – 26 August 2016,
Humboldt-Universität zu Berlin"

3pt functions and form factors

Z. Bajnok

Holographic QFT Group, Wigner Research Centre for Physics, Budapest

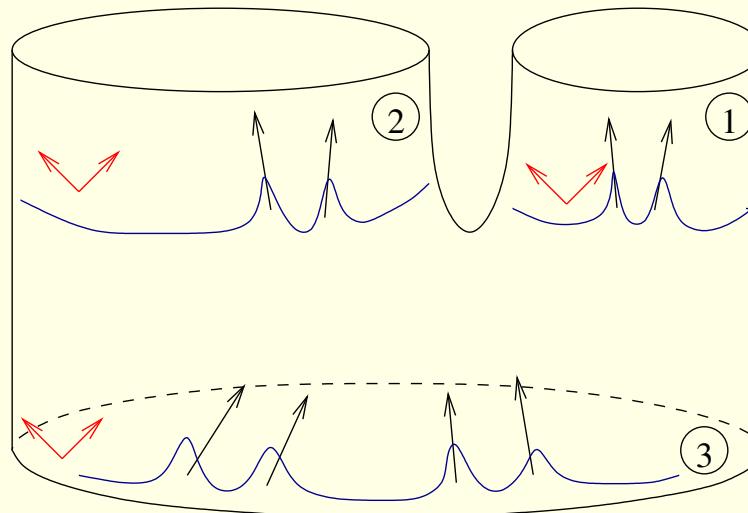
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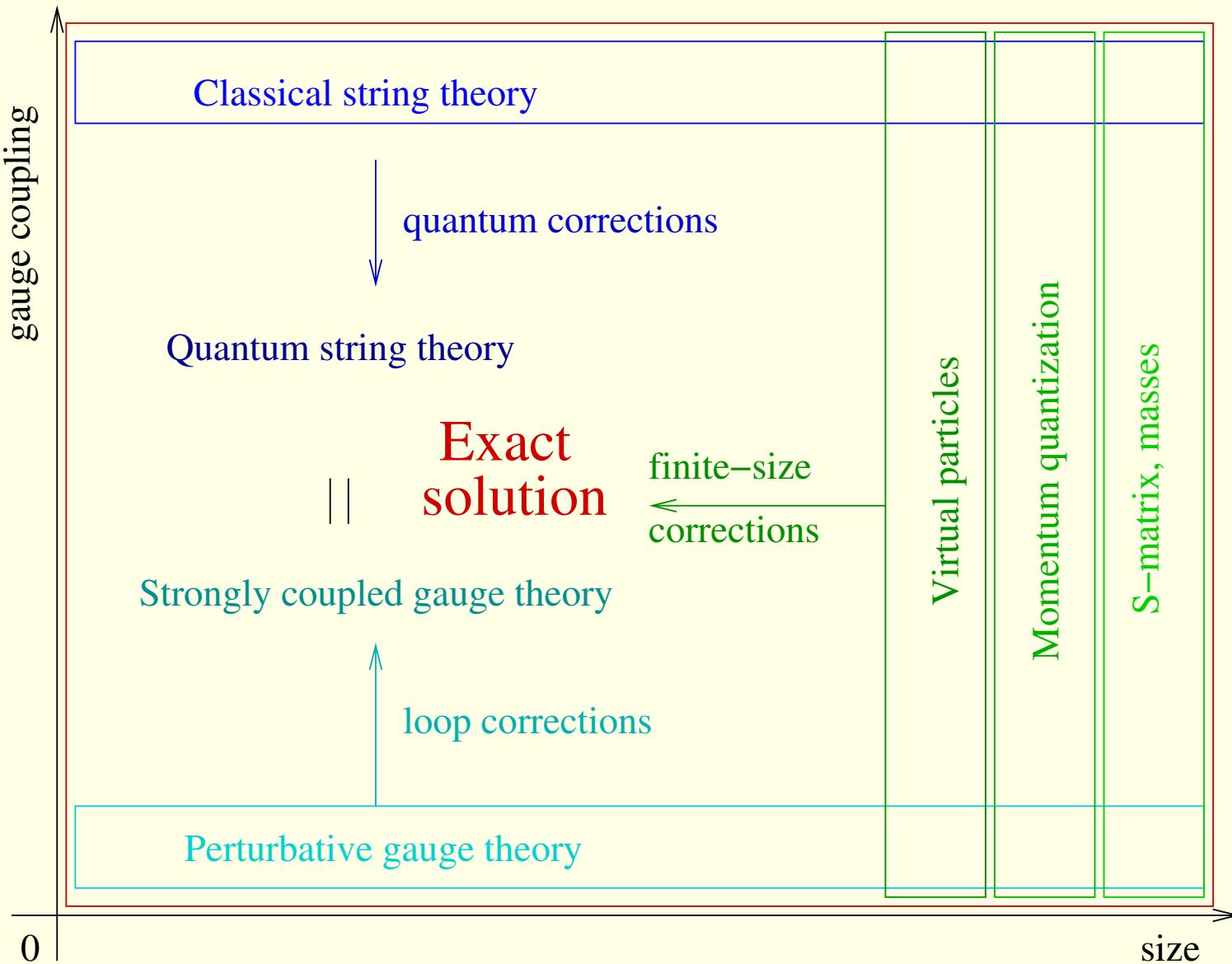
IIB strings on $AdS_5 \times S^5$	Integrability	$N = 4$ SYM
SFT vertex	Form factors	3pt functions



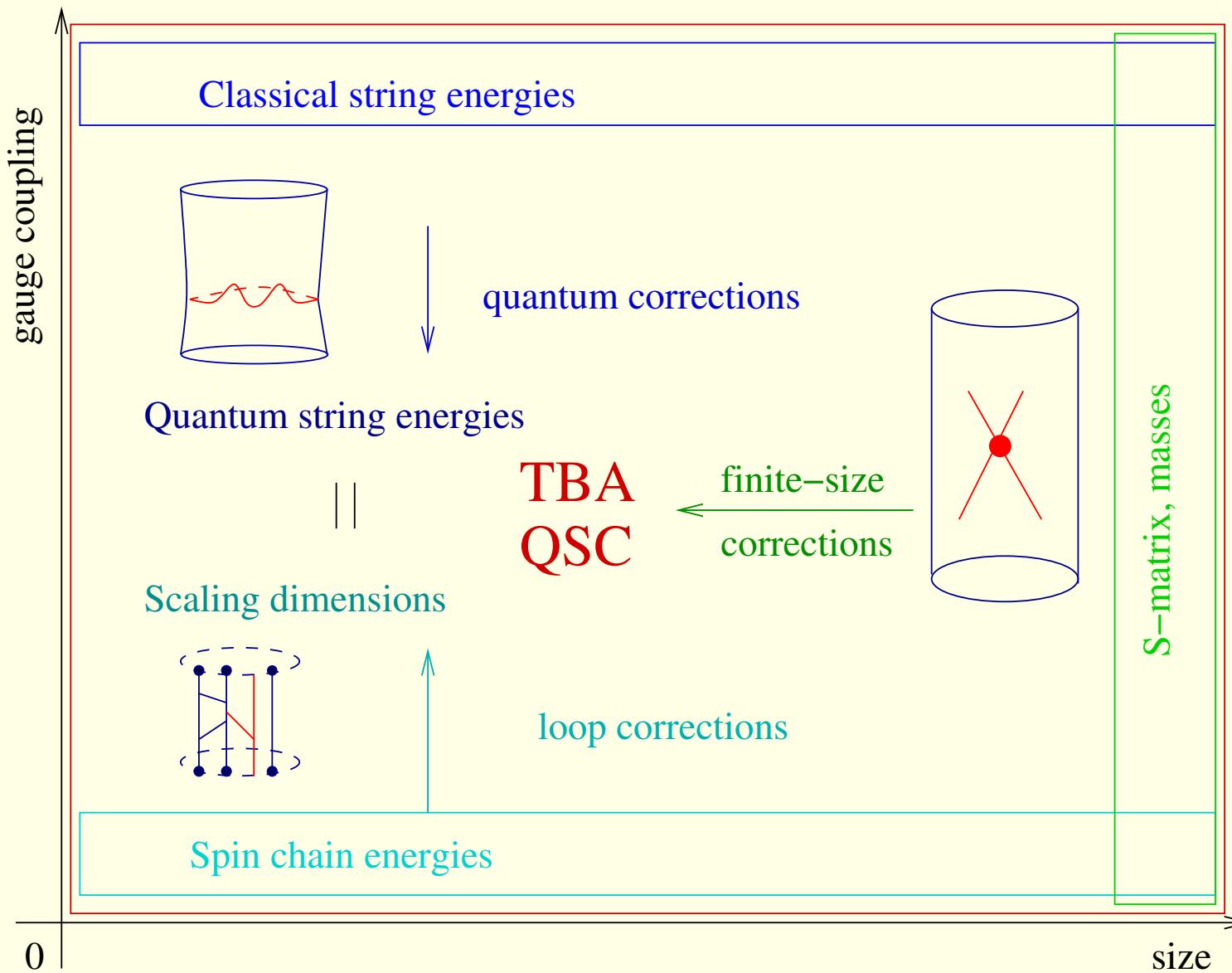
Same light-cone gauge fixing: integrable 2D QFT with particle like excitation. Amplitude \equiv string vertex \longleftrightarrow 3pt functions and $1/N$ corrections in dual gauge theory.

arXiv:1404.4556, 1501.04533, 1512.01471, 1607.02830: work done in collaboration with Romuald Janik and Andrzej Wereszczynski

Motivation:



Spectral problem: 2pt functions

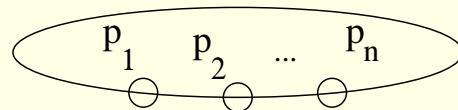


Motivation: insight from integrable models

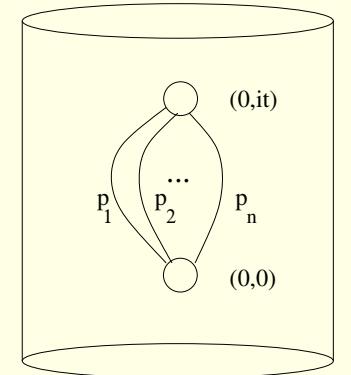
Motivation: insight from integrable models

The simplest interacting QFT in 1+1 D: $\mathcal{L} = \frac{1}{2}(\partial_t\varphi)^2 - \frac{1}{2}(\partial_x\varphi)^2 - \frac{m^2}{b^2}(\cosh b\varphi - 1)$

interesting observables: finite size spectrum,



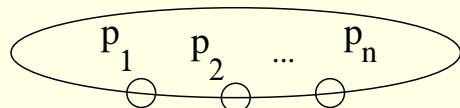
finite size correlators $L\langle 0|\mathcal{O}(it)\mathcal{O}(0)|0\rangle_L = \sum_n |L\langle 0|\mathcal{O}(0)|n\rangle_L|^2 e^{-E_n t}$



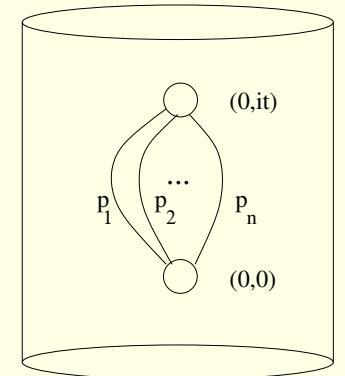
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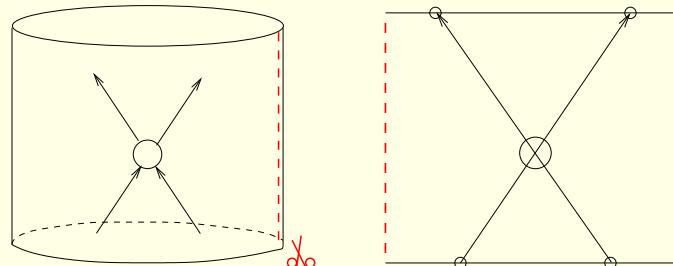


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Too difficult, instead

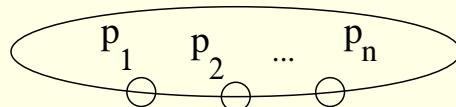
Infinite volume \rightarrow LSZ reduction formula



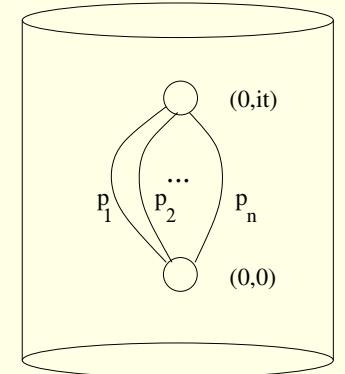
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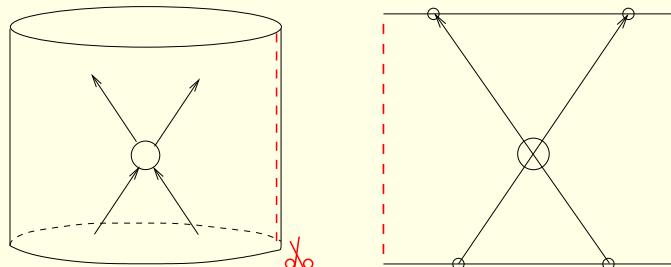


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$$\langle p'_1, p'_2 | \mathcal{O} | p_1, p_2 \rangle = \bar{\mathcal{D}}'_1 \bar{\mathcal{D}}'_2 \mathcal{D}_1 \mathcal{D}_2 \langle 0 | T(\mathcal{O} \varphi(1) \varphi(2) \varphi(3) \varphi(4)) | 0 \rangle$$

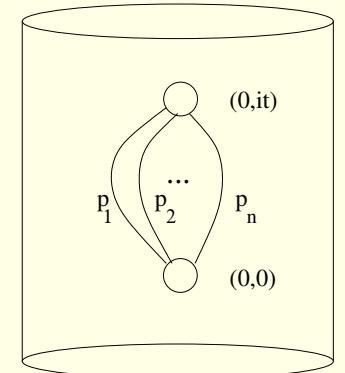
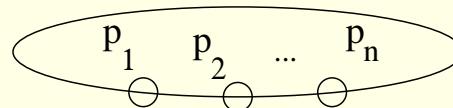
where $\mathcal{D}_j = i \int d^2x_j e^{ip_j x - i\omega_j t} \{ \partial_t^2 - \partial_x^2 + m^2 \}$ amputates a leg and puts it onshell

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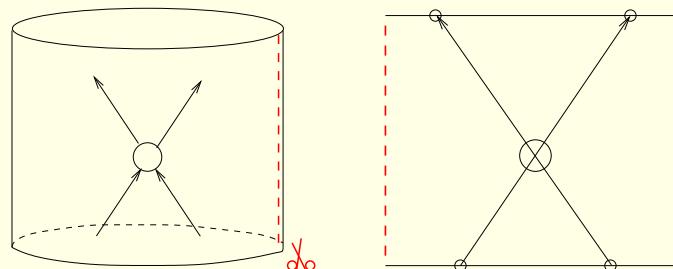
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Consequence: perturbative definition, calculational tool:

$$S(\theta) = 1 - \frac{1}{4}ib^2 \operatorname{csch}\theta - \frac{b^4 (\operatorname{csch}\theta (\pi \operatorname{csch}\theta - i))}{32\pi} + \frac{ib^6 \operatorname{csch}\theta (\pi \operatorname{csch}\theta - i)^2}{256\pi^2} + O(b^8)$$

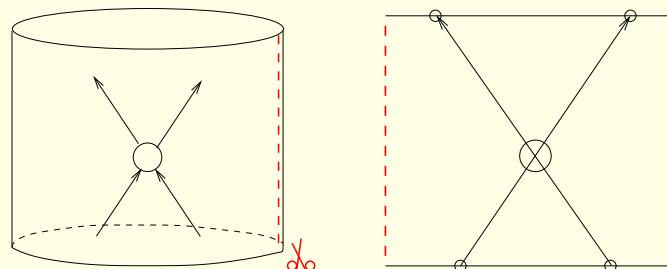
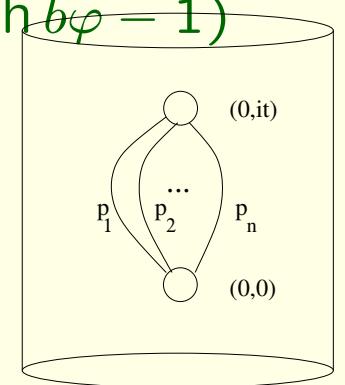
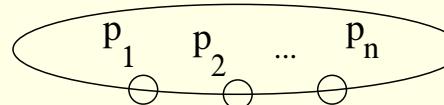
Mandelstam variable $s = 4m^2 \cosh^2 \frac{\theta}{2}$ where $\theta = \theta_1 - \theta_2$ rapidity: $p_i = m \sinh \theta_i$

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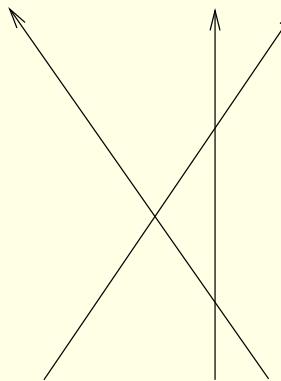
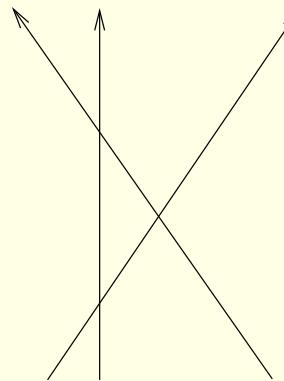
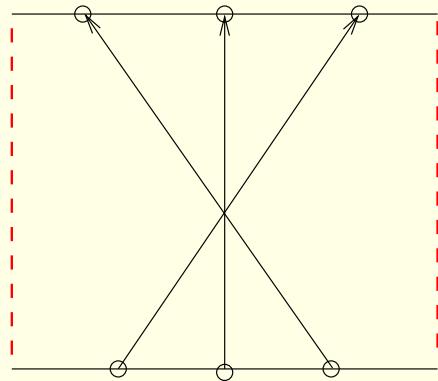
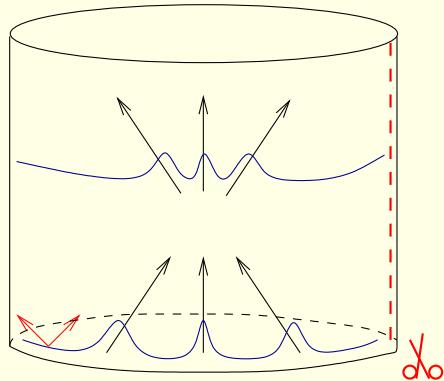
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control over analytical properties: unitarity, crossing $S(\theta) = S(-\theta)^{-1} = S(i\pi - \theta)$

Motivation: S-matrix bootstrap

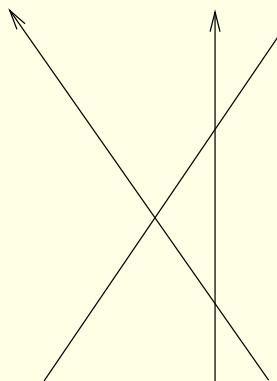
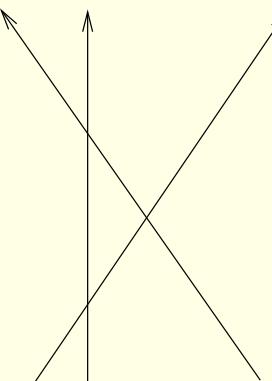
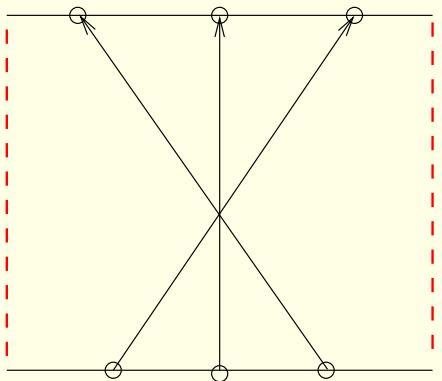
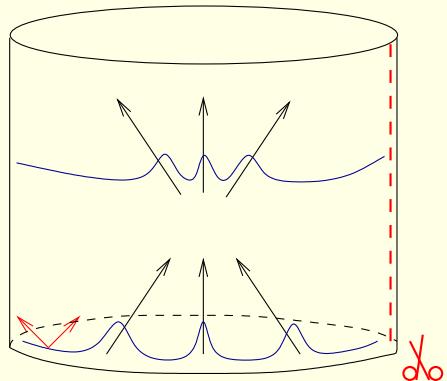
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S-matrix bootstrap: fundamental object is the two particle S-matrix [Zamolodchikov² '79,Dorey]

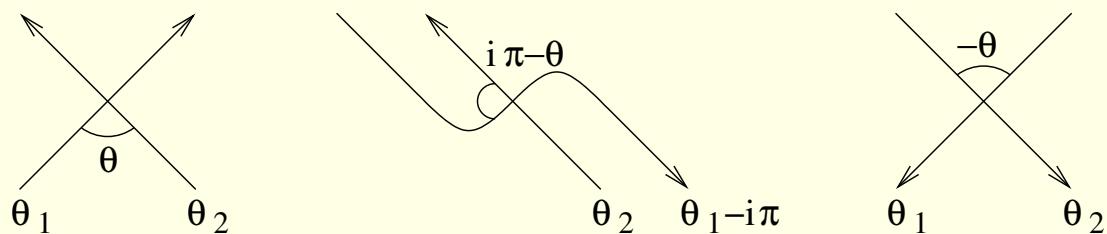


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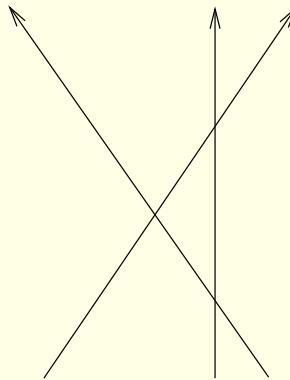
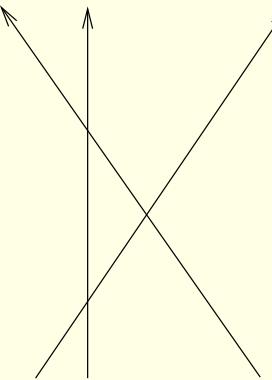
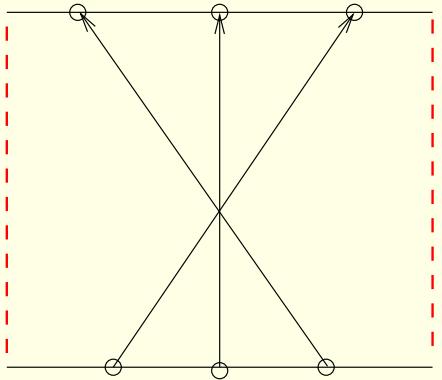
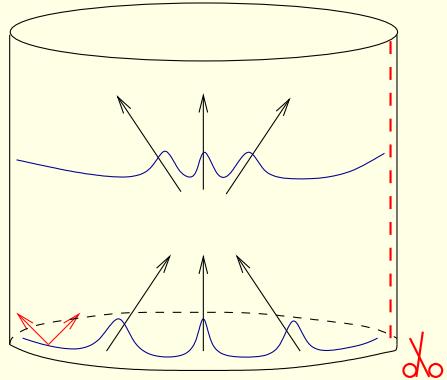
Infinite volume \rightarrow crossing symmetry, $\theta \rightarrow i\pi - \theta$ in rapidity $(E(\theta), p(\theta)) = m(\cosh \theta, \sinh \theta)$



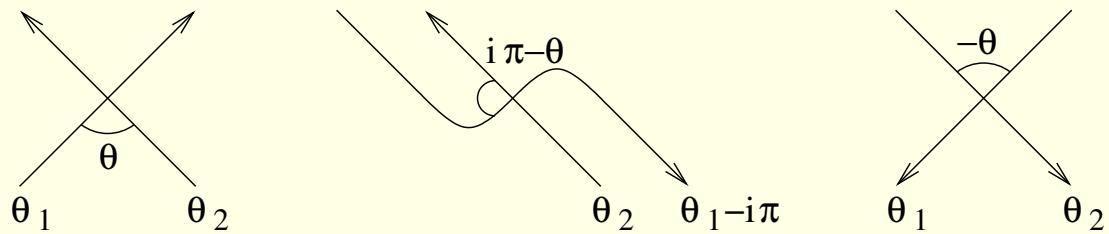
$$S(\theta_1 - \theta_2) = S(\theta) = S(i\pi - \theta) = S(-\theta)^{-1} :$$

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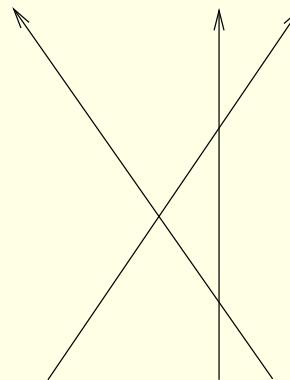
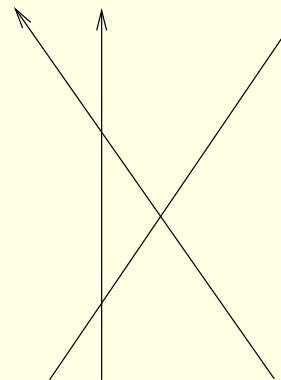
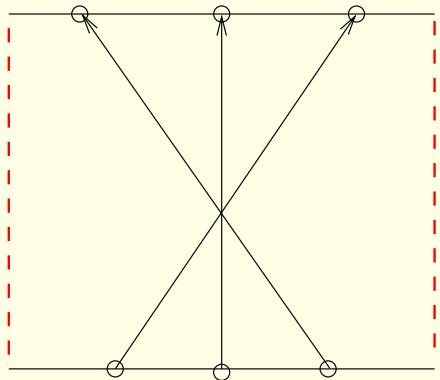
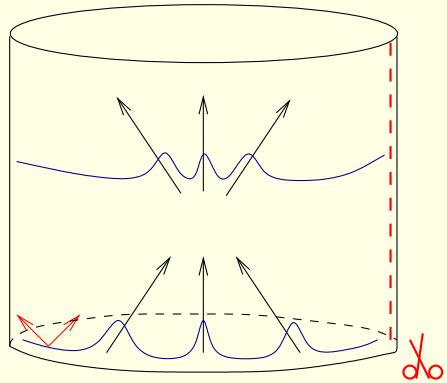
Simple solution:
sinh-Gordon
 $S(\theta) = \frac{\sinh \theta - ia}{\sinh \theta + ia}$

$$a = \frac{\pi b^2}{8\pi + b^2}$$

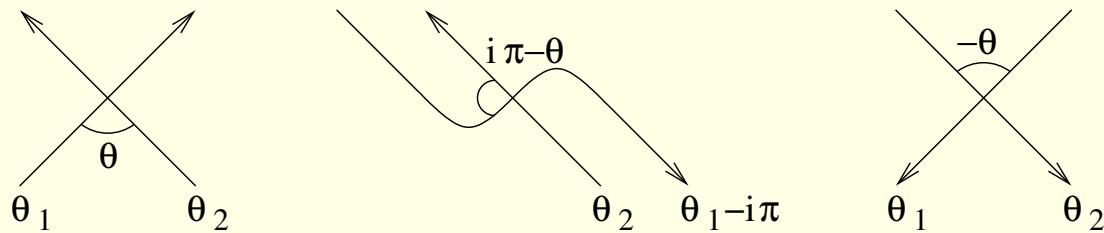
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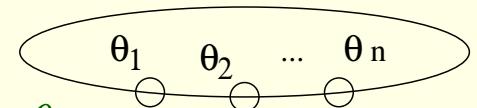
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Finite volume spectrum [Bethe-Yang] upto $O(e^{-mL})$ i.e. polynomial in L^{-1} :

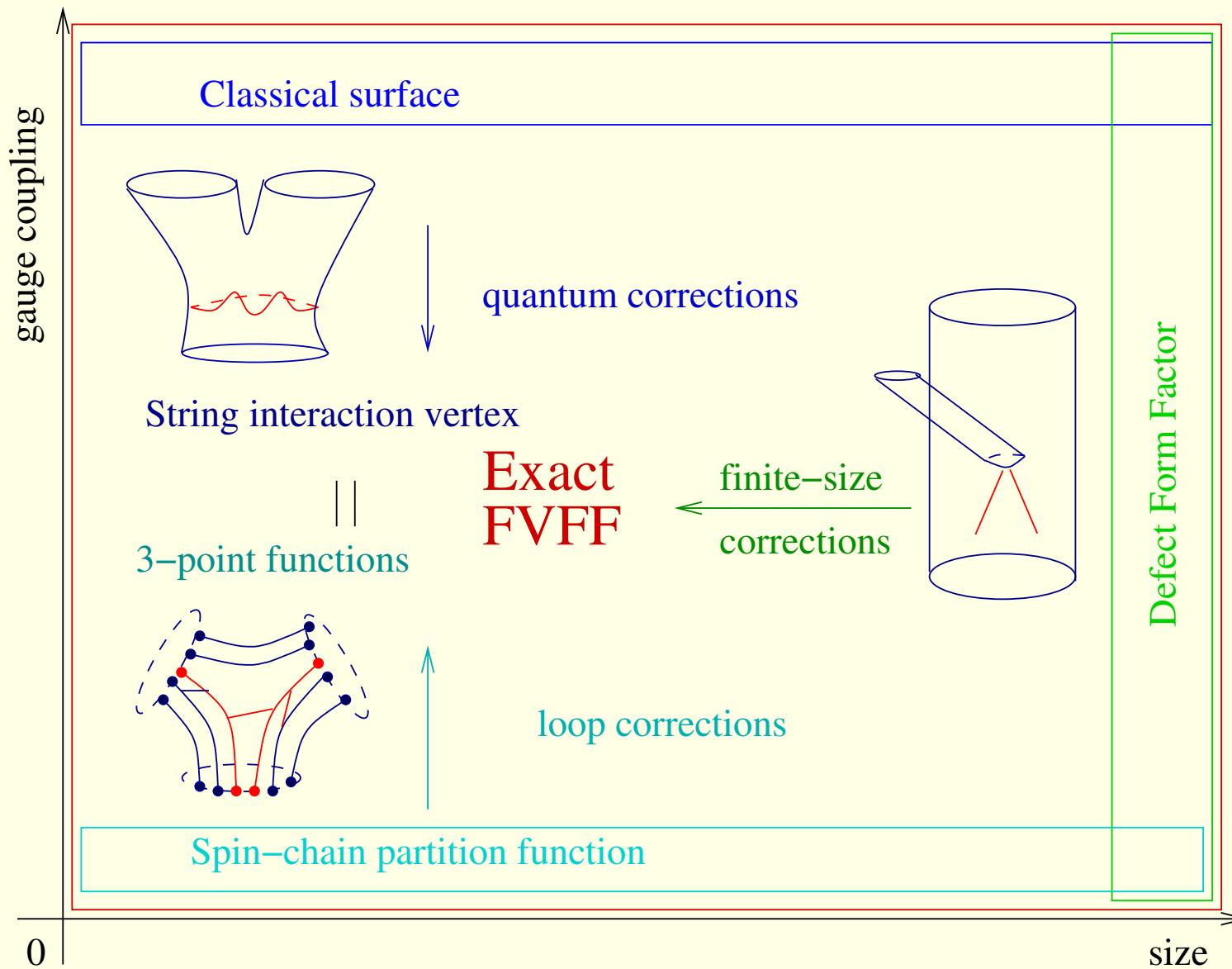
$$e^{i\Phi_1} = e^{ip_1 L} S(\theta_1 - \theta_2) \dots S(\theta_1 - \theta_n) = 1 \quad E_n(L) = \sum_i m \cosh \theta_i$$

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String interaction, 3pt functions

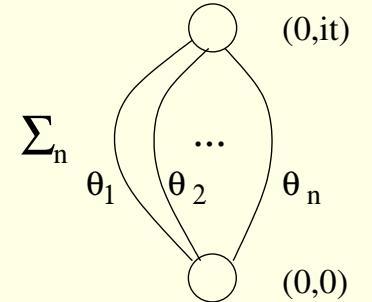


Form factor bootstrap

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Correlation functions: [Smirnov, Karowszki] $\langle 0 | \mathcal{O}(it) \mathcal{O}(0) | 0 \rangle =$

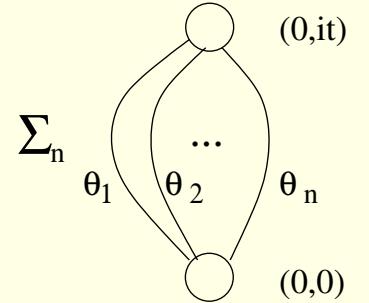
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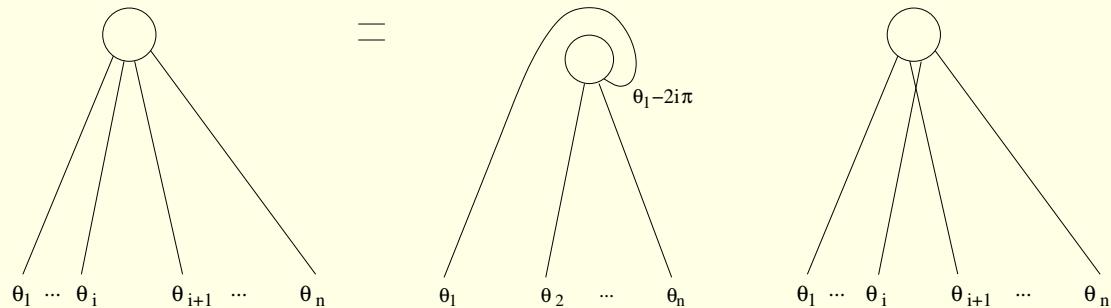
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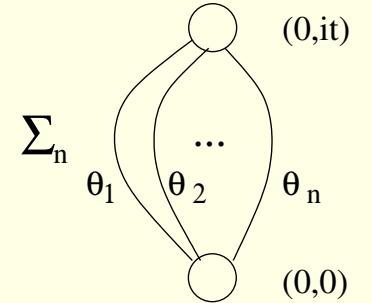


$$\langle 0 | \mathcal{O} | \theta_1, \dots, \theta_n \rangle = \langle 0 | \mathcal{O} | \theta_2, \dots, \theta_n, \theta_1 - 2i\pi \rangle = S(\theta_i - \theta_{i+1}) \langle 0 | \mathcal{O} | \dots, \theta_{i+1}, \theta_i, \dots \rangle$$

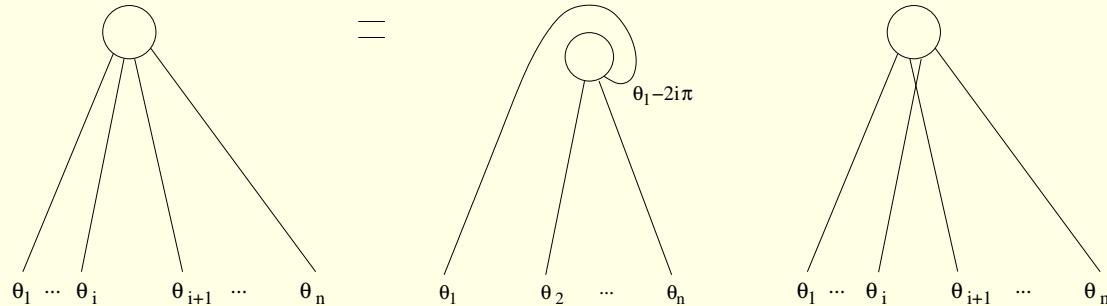
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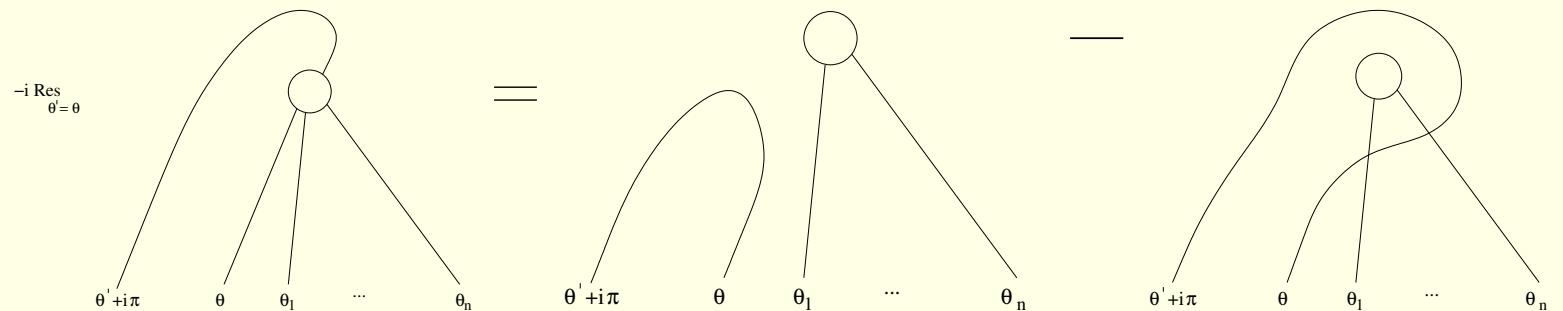


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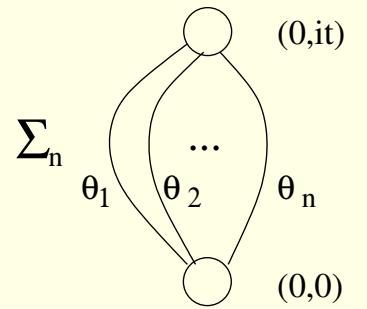
Singularity structure



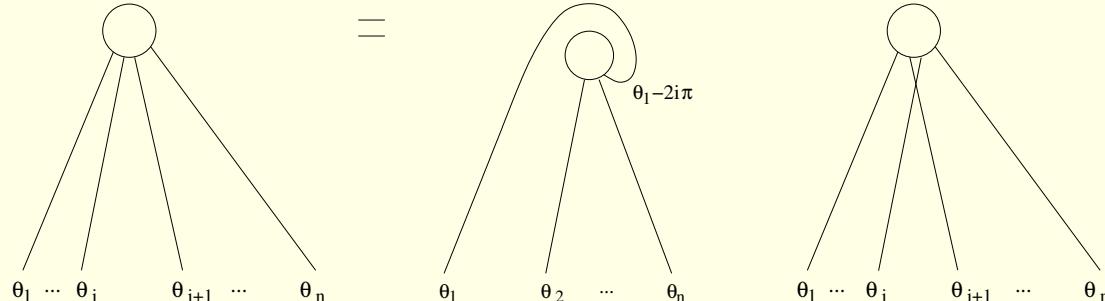
$$-i \text{Res}_{\theta'=\theta} \langle 0 | \mathcal{O} | \theta' + i\pi, \theta, \theta_1, \dots, \theta_n \rangle = (1 - \prod_i S(\theta - \theta_i)) \langle 0 | \mathcal{O} | \theta_1, \dots, \theta_n \rangle$$

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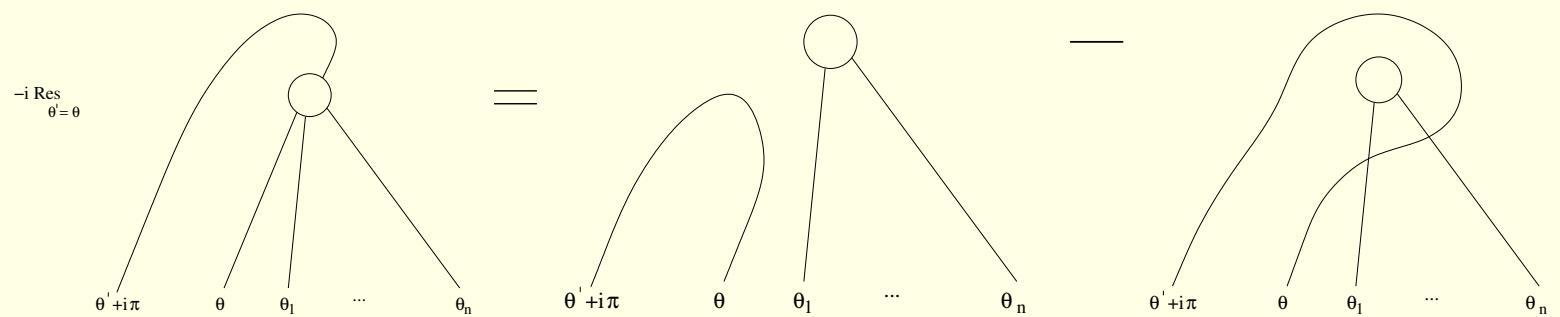


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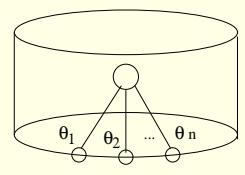
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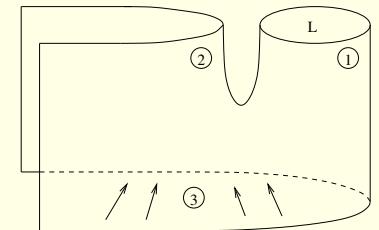
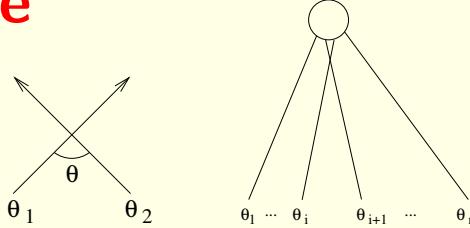
Solution for sinh-Gordon: $\langle 0|\mathcal{O}|\theta_1, \theta_2 \rangle = e^{(D+D^{-1})^{-1} \log S}; Df(\theta) = f(\theta + i\pi)$

Finite volume form factors: polynomial in L^{-1} : $\langle 0|\mathcal{O}|\theta_1, \dots, \theta_n \rangle_L = \frac{\langle 0|\mathcal{O}|\theta_1, \dots, \theta_n \rangle}{\sqrt{\det[\frac{\partial \Phi_i}{\partial \theta_j}]}}$

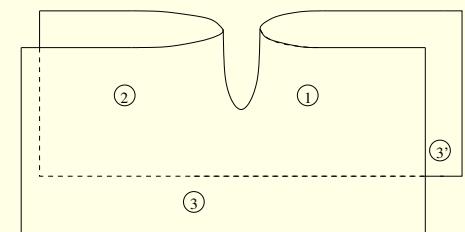
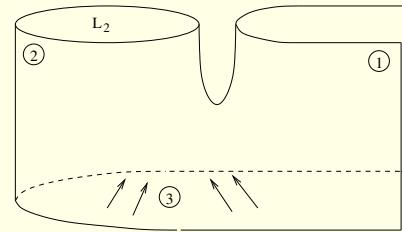


Outline

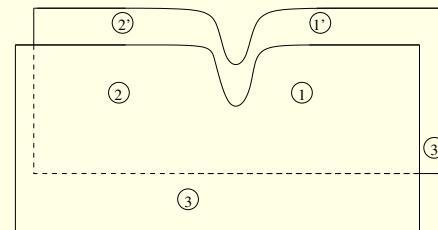
Motivation: the S-matrix and the FF bootstrap



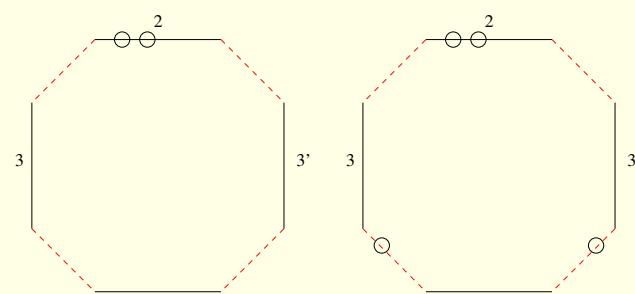
Cutting the stringvertex: decompactifying 2 & 3: Form factor axioms



Decompacting all volumes: Octagon axioms



Cutting one more: hexagon axioms



PP wave limit of AdS/CFT: Solving the FF axioms

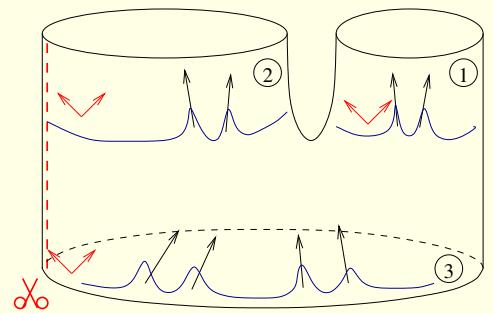
Summing up the octagon

Conclusion

Decompactification limit of the string vertex

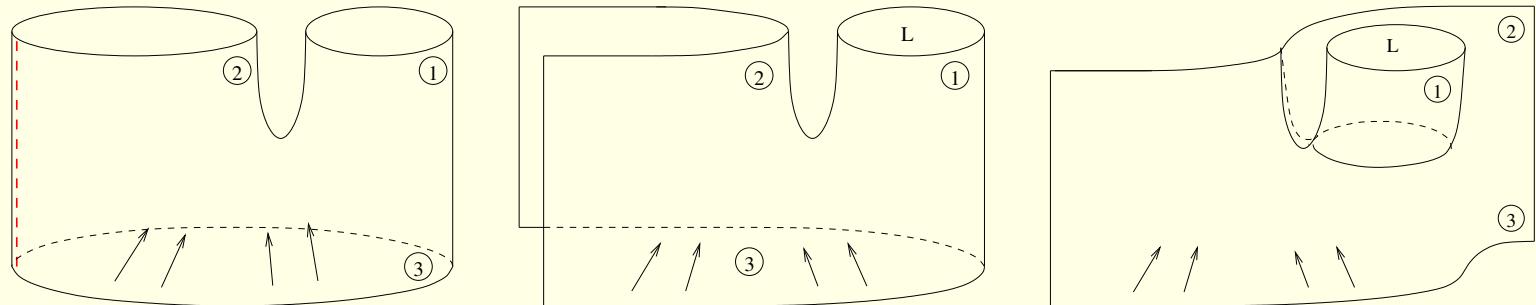
Decompactification limit of the string vertex

Decompactify string 2 & 3:



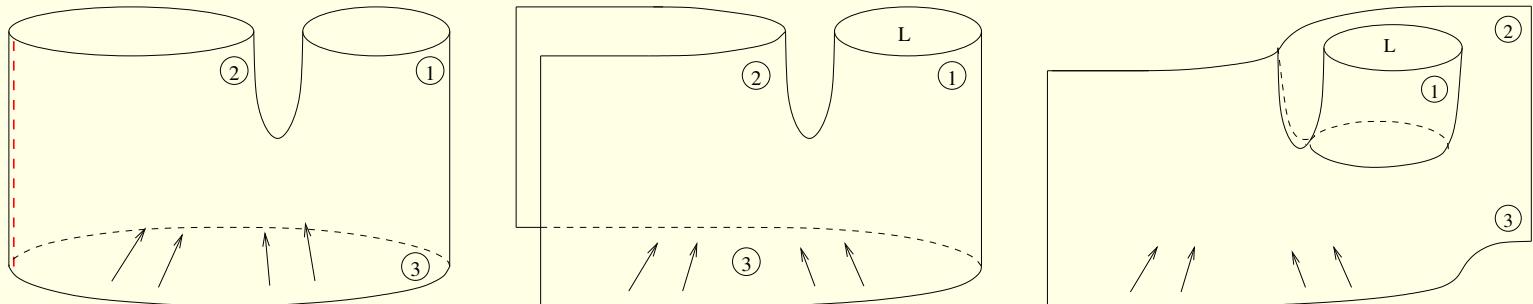
Decompactification limit of the string vertex

Decompactify
string 2 & 3:

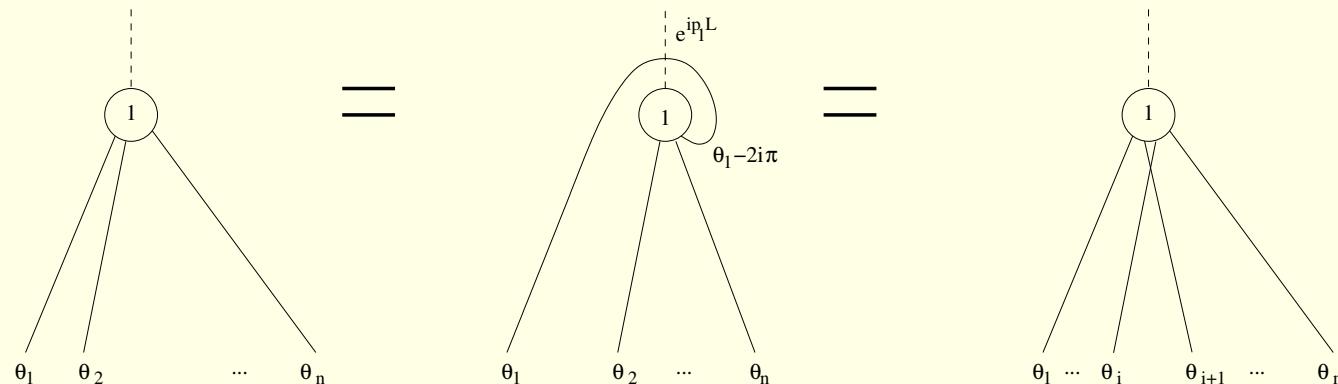


Decompactification limit of the string vertex

Decompactify
string 2 & 3:



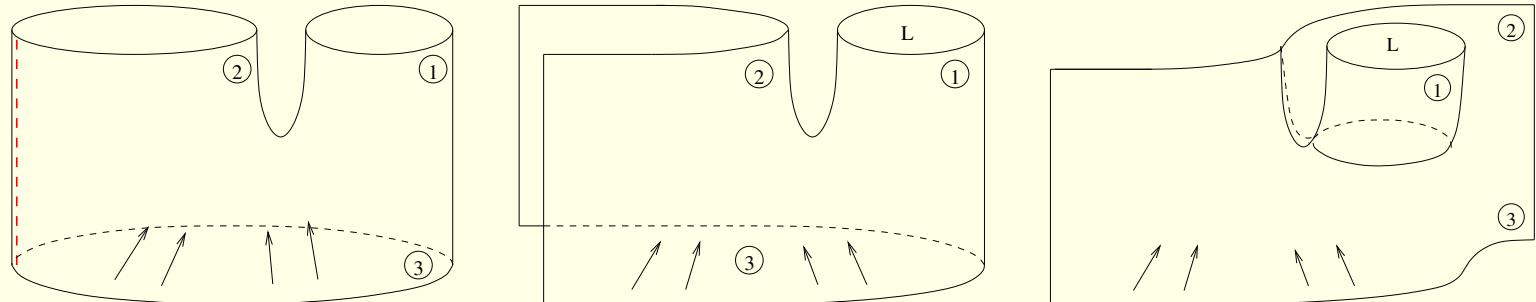
Form factor equations:



$$N_L(\theta_1, \dots, \theta_n) = e^{-ip_1 L} N_L(\theta_2, \dots, \theta_n, \theta_1 - 2i\pi) = S(\theta_i - \theta_{i+1}) N_L(\dots, \theta_{i+1}, \theta_i, \dots)$$

Decompactification limit of the string vertex

Decompactify
string 2 & 3:

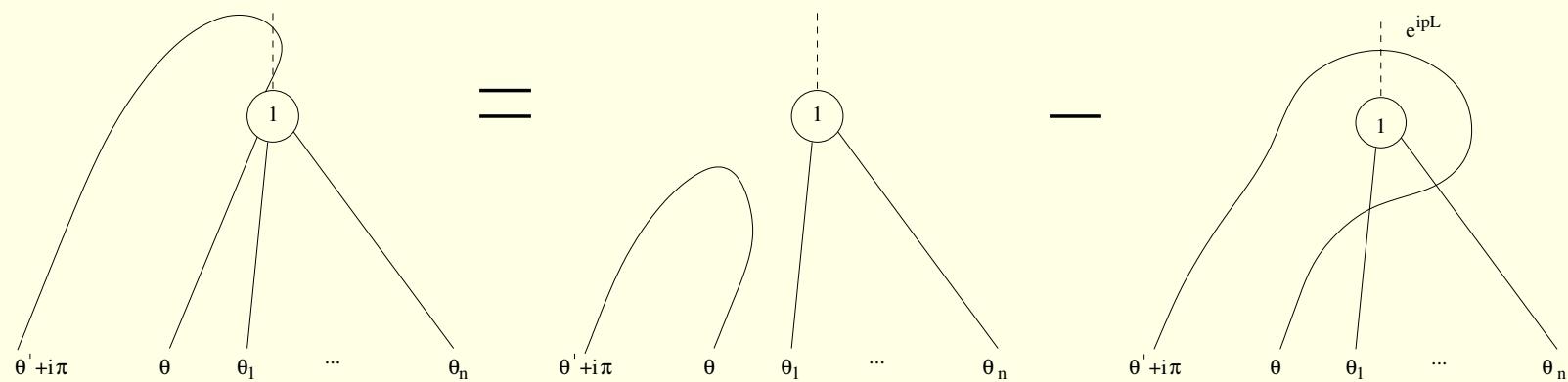


Form factor equations:

$$\begin{array}{c}
 \text{Diagram 1: } \text{A circle labeled 1 connected by dashed lines to } \theta_1, \theta_2, \dots, \theta_n \\
 = \\
 \text{Diagram 2: } \text{A circle labeled 1 connected by dashed lines to } \theta_1, \theta_2, \dots, \theta_n, \text{ with a loop above it labeled } e^{ip_1 L} \text{ and } \theta_1 - 2i\pi \\
 = \\
 \text{Diagram 3: } \text{A circle labeled 1 connected by dashed lines to } \theta_1, \dots, \theta_i, \theta_{i+1}, \dots, \theta_n
 \end{array}$$

$$N_L(\theta_1, \dots, \theta_n) = e^{-ip_1 L} N_L(\theta_2, \dots, \theta_n, \theta_1 - 2i\pi) = S(\theta_i - \theta_{i+1}) N_L(\dots, \theta_{i+1}, \theta_i, \dots)$$

Singularity structure



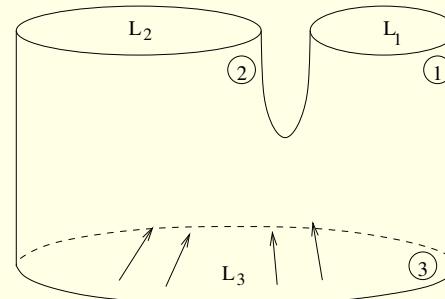
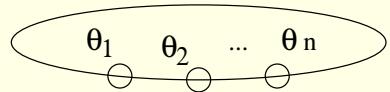
$$-i \text{Res}_{\theta'=\theta} N_L(\theta' + i\pi, \theta, \theta_1, \dots, \theta_n) = (1 - e^{ipL} \prod_i S(\theta - \theta_i)) N_L(\theta_1, \dots, \theta_n)$$

Other decompactification: the bootstrap program

Other decompactification: the bootstrap program

The full large volume amplitude $O(e^{-mL})$

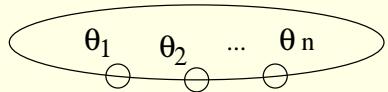
$$e^{ip_1 L} S(\theta_1 - \theta_2) \dots S(\theta_1 - \theta_n) = 1$$



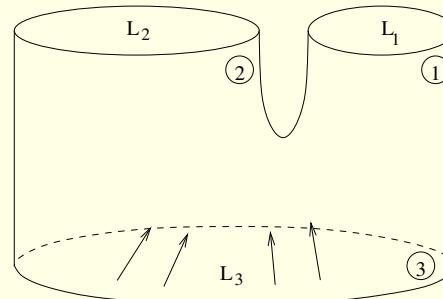
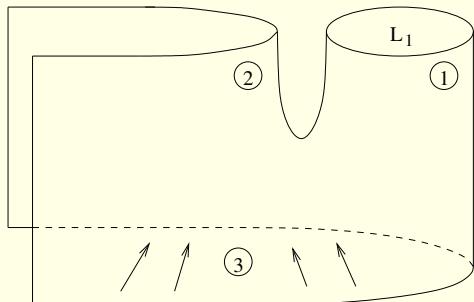
Other decompactification: the bootstrap program

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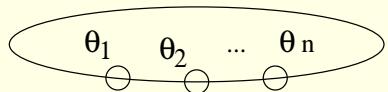
Decompactify
string 2 & 3:
 $N_{L_1}(\theta_1, \dots, \theta_n)$



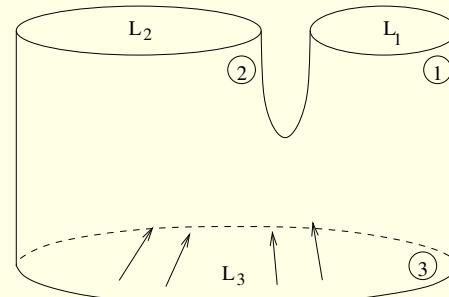
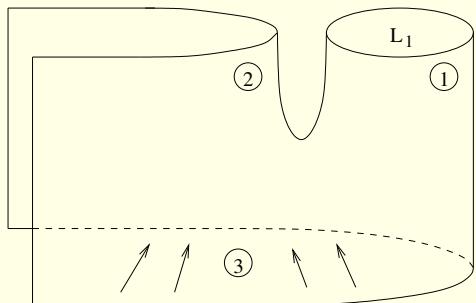
Other decompactification: the bootstrap program

The full large volume amplitude $O(e^{-mL})$

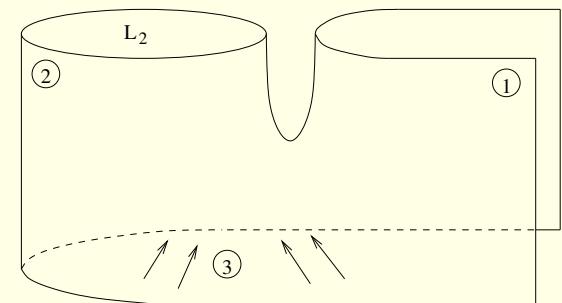
$$e^{ip_1 L} S(\theta_1 - \theta_2) \dots S(\theta_1 - \theta_n) = 1$$



Decompactify
string 2 & 3:
 $N_{L_1}(\theta_1, \dots, \theta_n)$



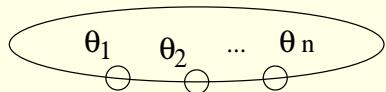
Decompactify
string 1 & 3:
 $N_{L_2}(\theta_1, \dots, \theta_n)$



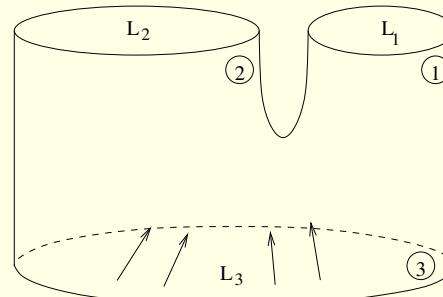
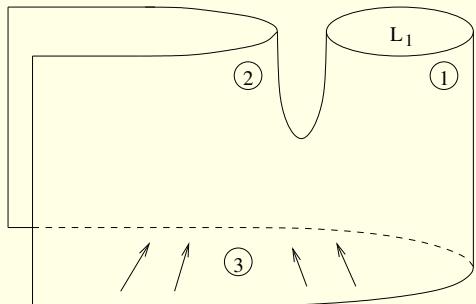
Other decompactification: the bootstrap program

The full large volume amplitude $O(e^{-mL})$

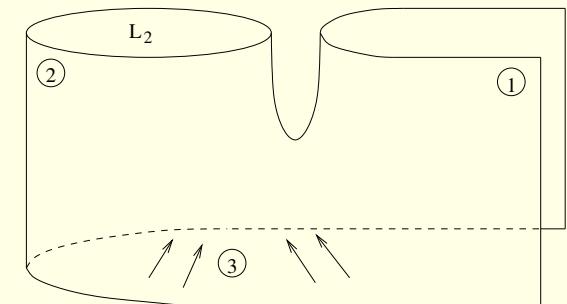
$$e^{ip_1 L} S(\theta_1 - \theta_2) \dots S(\theta_1 - \theta_n) = 1$$



Decompactify
string 2 & 3:
 $N_{L_1}(\theta_1, \dots, \theta_n)$



Decompactify
string 1 & 3:
 $N_{L_2}(\theta_1, \dots, \theta_n)$



Finite (large volume) and infinite volume amplitudes are the same (upto normalization).

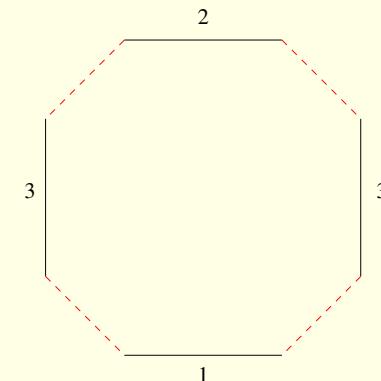
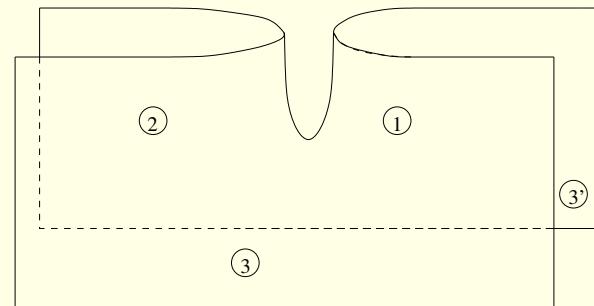
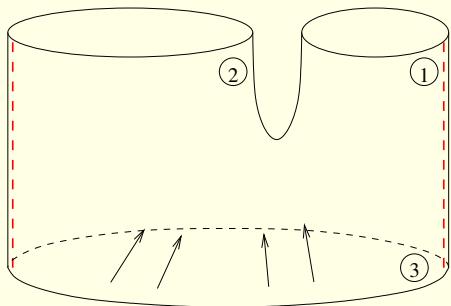
Find the relevant solutions by matching the two in the large L_1, L_2 limit:

$$N_{L_1}(\theta_1, \dots, \theta_n) \propto N_{L_1, L_2}(\theta_1, \dots, \theta_n) \propto N_{L_2}(\theta_1, \dots, \theta_n)$$

Decompactifying all volumes $L_1 = \infty$: octagon

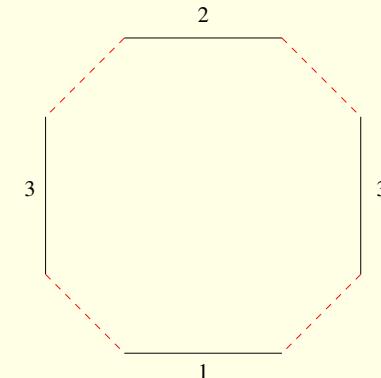
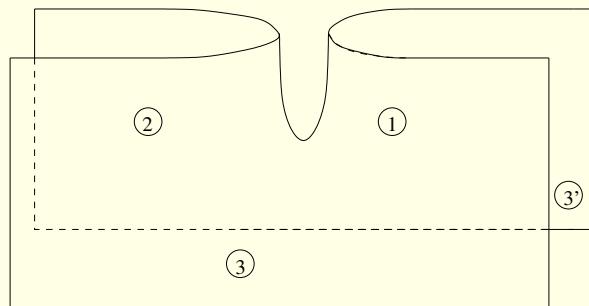
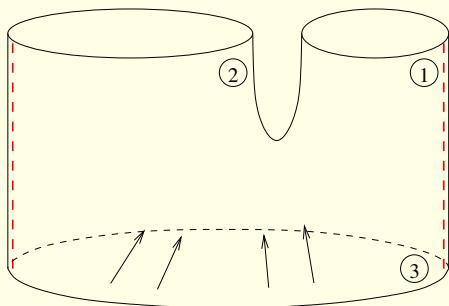
Decompactifying all volumes $L_1 = \infty$: octagon

Decompactify all volumes

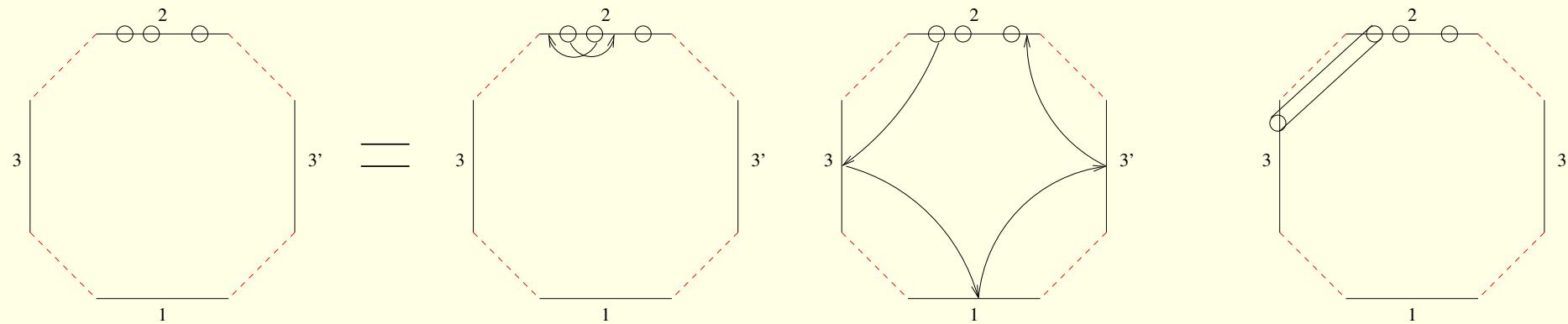


Decompactifying all volumes $L_1 = \infty$: octagon

Decompactify all volumes



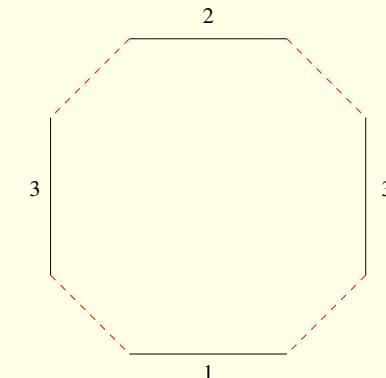
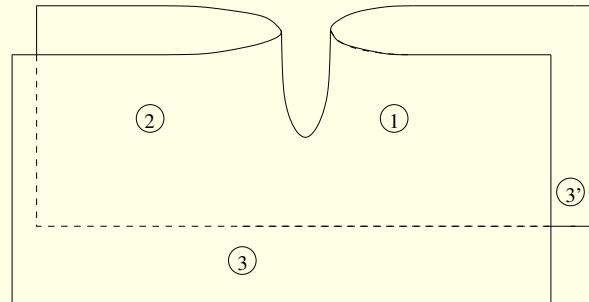
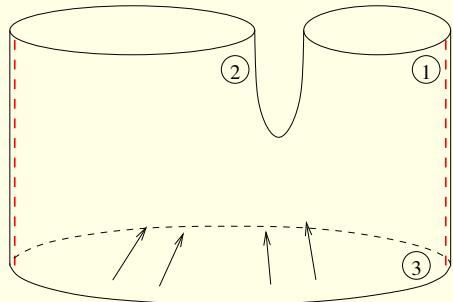
Octagon axioms:



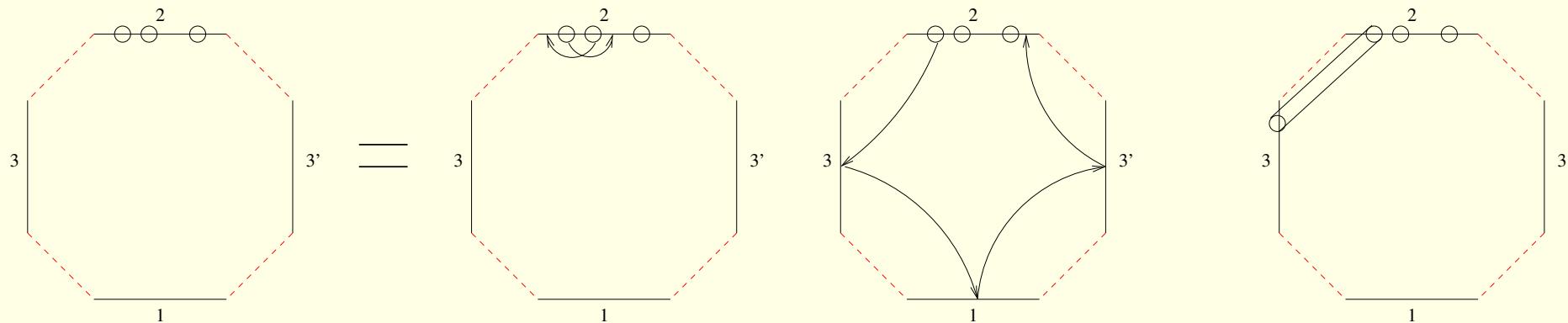
$$O(\theta_i, \dots, \theta_n) = S(\theta_i, \theta_{i+1}) O(\dots, \theta_{i+1}, \theta_i, \dots) = O(\theta_2, \dots, \theta_n, \theta_1 - 4i\pi)$$

Decompactifying all volumes $L_1 = \infty$: octagon

Decompactify all volumes



Octagon axioms:



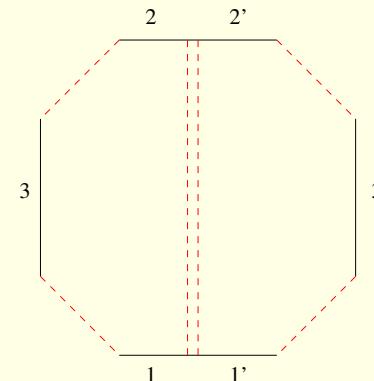
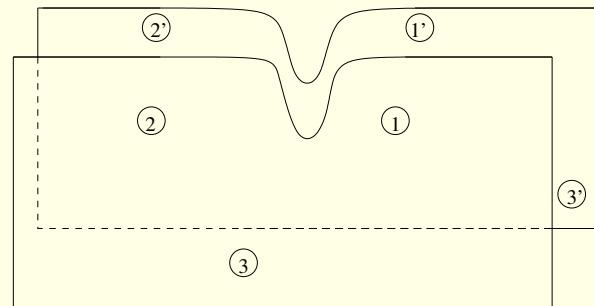
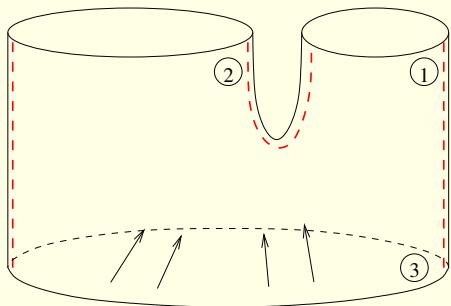
$$O(\theta_i, \dots, \theta_n) = S(\theta_i, \theta_{i+1}) O(\dots, \theta_{i+1}, \theta_i, \dots) = O(\theta_2, \dots, \theta_n, \theta_1 - 4i\pi)$$

$$\text{Kinematical singularity } -i\text{Res}_{\theta'=\theta} O(\theta' + i\pi, \theta, \theta_1, \dots, \theta_n) = O(\theta_1, \dots, \theta_n)$$

Cutting one more: two hexagons

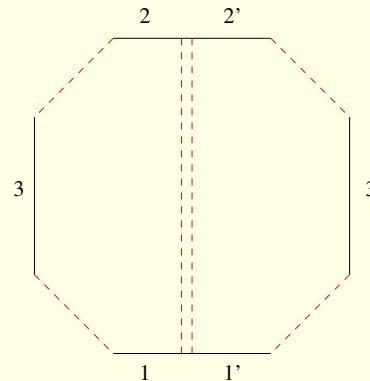
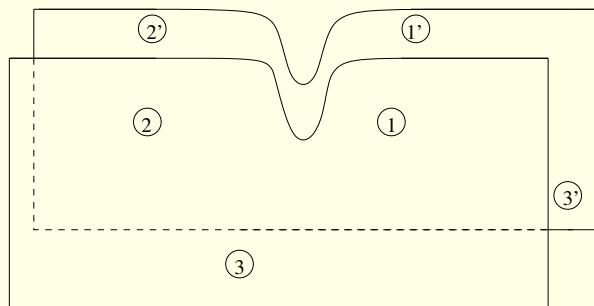
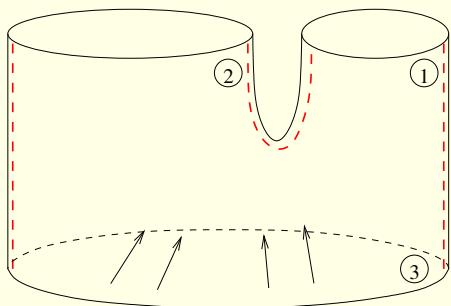
Cutting one more: two hexagons

Decompatify all volumes

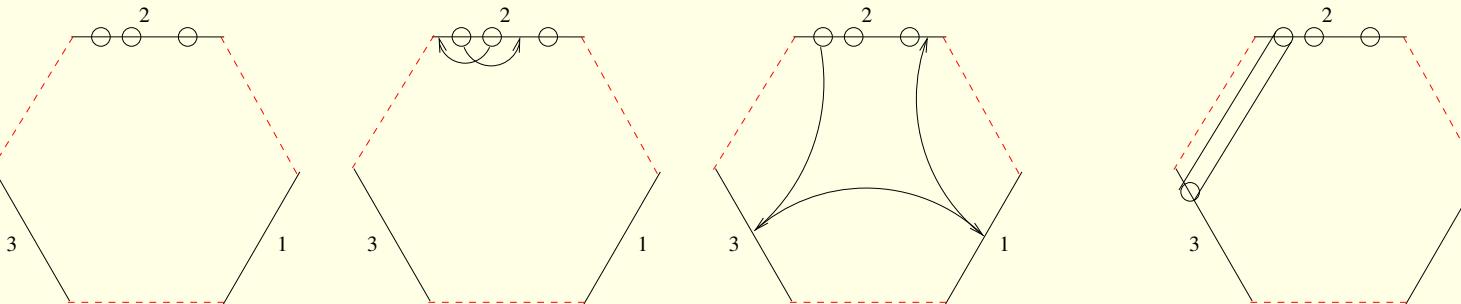


Cutting one more: two hexagons

Decompatify all volumes



Hexagon axioms: [Basso,Komatsu,Vieira '15]

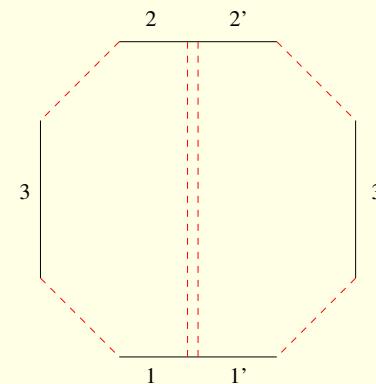
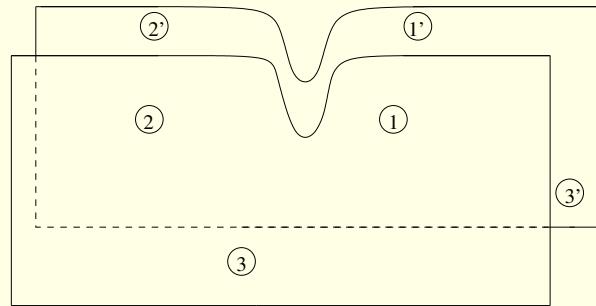
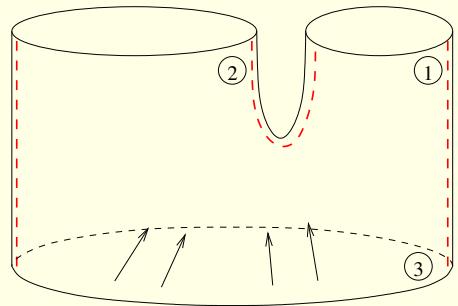


$$h(\theta_i, \dots, \theta_n) = S(\theta_i, \theta_{i+1}) h(\dots, \theta_{i+1}, \theta_i, \dots) = h(\theta_2, \dots, \theta_n, \theta_1 - 3i\pi)$$

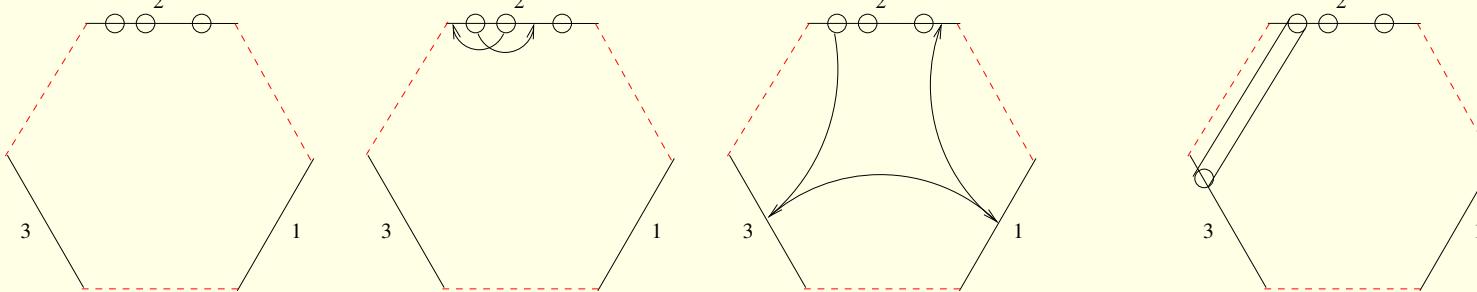
$$\text{Kinematical singularity } -i\text{Res}_{\theta'=\theta} h(\theta' + i\pi, \theta, \theta_1, \dots, \theta_n) = h(\theta_1, \dots, \theta_n)$$

Cutting one more: two hexagons

Decompactify all volumes



Hexagon axioms: [Basso,Komatsu,Vieira '15]



$$h(\theta_i, \dots, \theta_n) = S(\theta_i, \theta_{i+1}) h(\dots, \theta_{i+1}, \theta_i, \dots) = h(\theta_2, \dots, \theta_n, \theta_1 - 3i\pi)$$

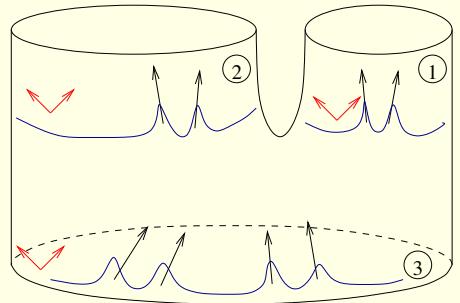
$$\text{Kinematical singularity } -i\text{Res}_{\theta'=\theta} h(\theta' + i\pi, \theta, \theta_1, \dots, \theta_n) = h(\theta_1, \dots, \theta_n)$$

$$\text{Complete solution: } h(\theta_1, \theta_2) \propto \sigma(\theta_1, \theta_2) S_{\text{Beisert}}(\theta_1, \theta_2)$$

Comparision of the different approaches

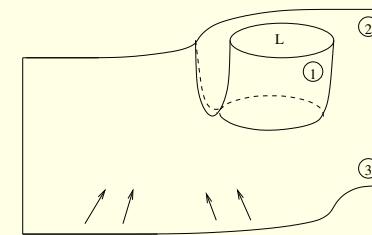
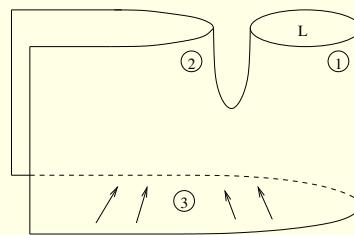
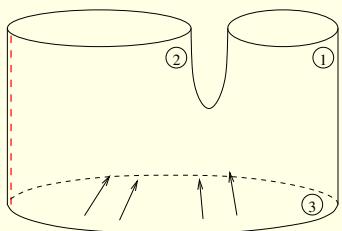
Comparision of the different approaches

Ultimate goal:



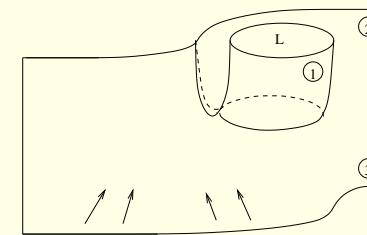
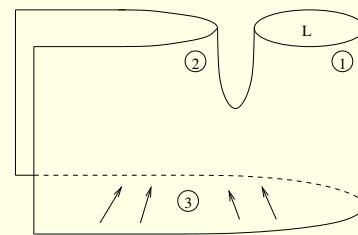
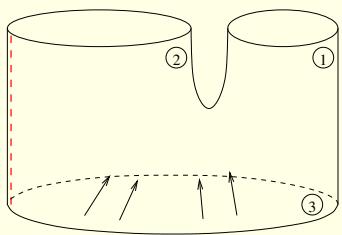
Comparision of the different approaches

1 cut:
nonlocal
form factors

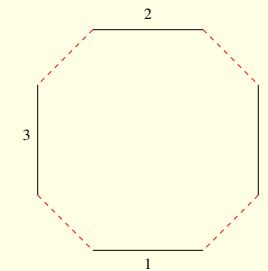
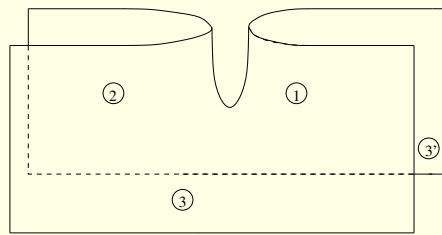
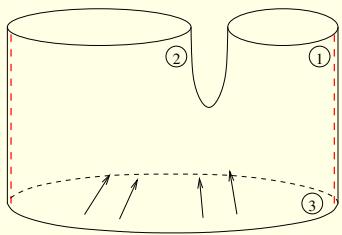


Comparision of the different approaches

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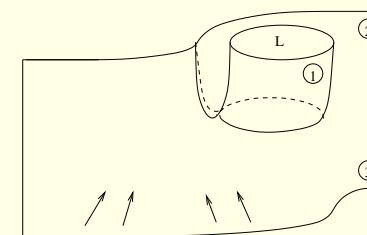
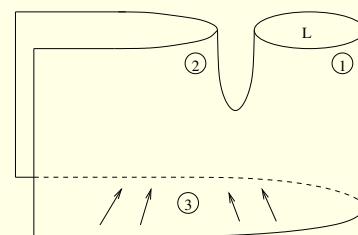
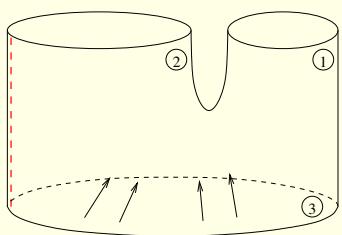


2 cuts
octagon

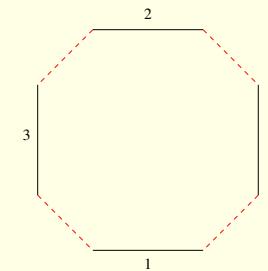
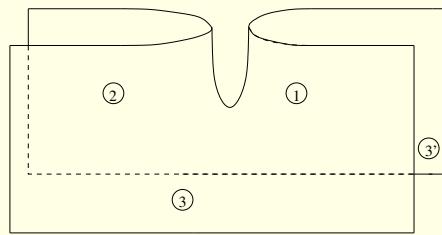
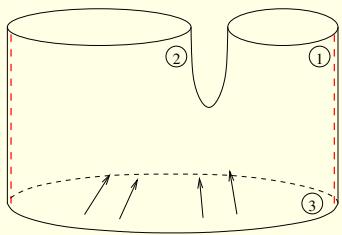


Comparision of the different approaches

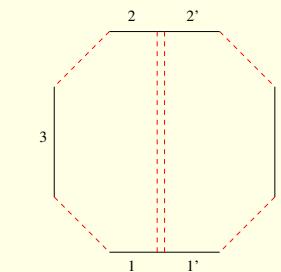
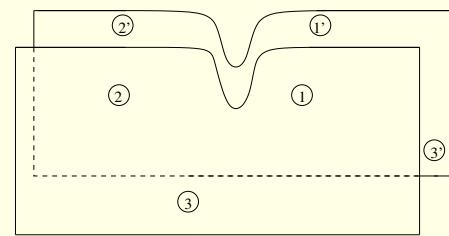
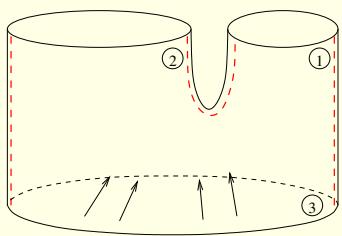
1 cut:
nonlocal
form factors



2 cuts
octagon

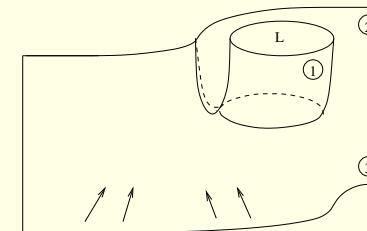
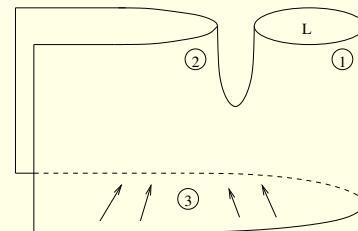
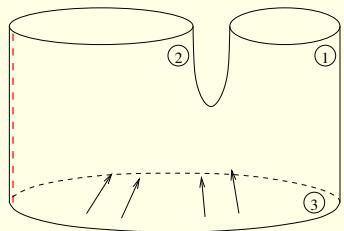


3 cuts
hexagon

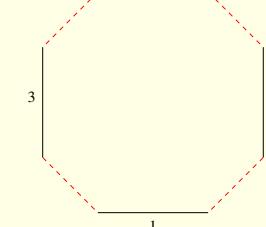
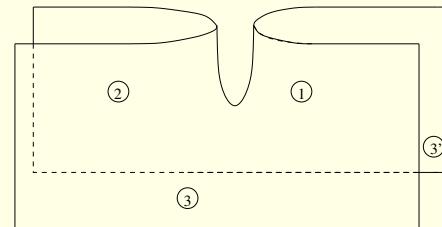
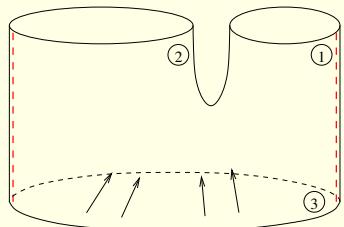


Comparision of the different approaches

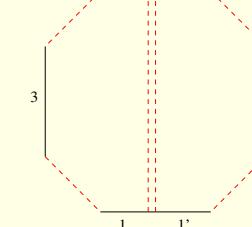
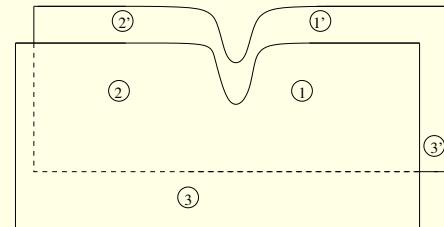
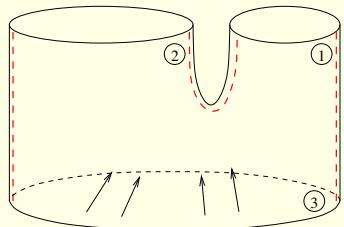
1 cut:
nonlocal
form factors



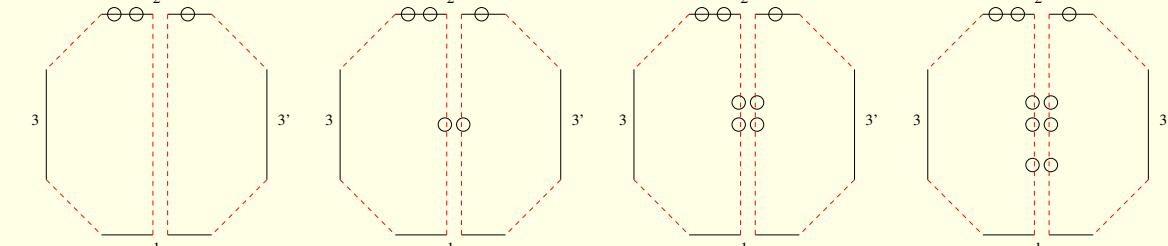
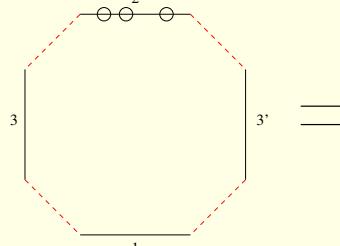
2 cuts
octagon



3 cuts
hexagon



sewing back



$$O(\theta_1, \theta_2, \theta_3) = h(\theta_1, \theta_2)h(\theta_3) + \dots + \int \frac{du}{2\pi} \mu(u) h(\theta_1, \theta_2, u - i\frac{\pi}{2}) h(u + i\frac{\pi}{2}, \theta_3) e^{-E(u)l} +$$

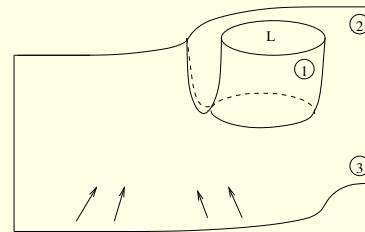
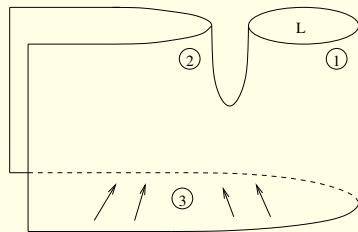
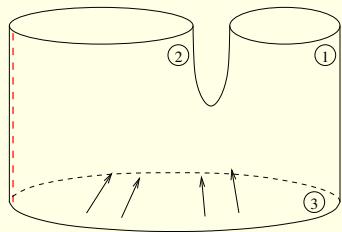
Perturbative checks: [Eden & Sfondrini] [Basso et. al]

HHH: [Jiang, Komatsu, Kostov, Serban]

Free massive boson: pp-wave limit of strings on $AdS_5 \times S^5$

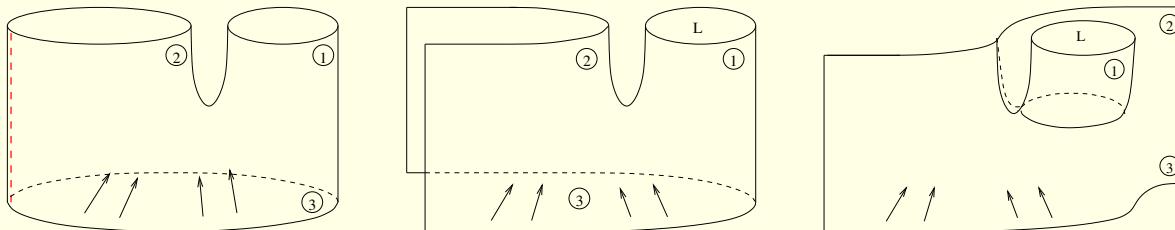
Free massive boson: pp-wave limit of strings on $AdS_5 \times S^5$

1 cut:
nonlocal
form factors



Free massive boson: pp-wave limit of strings on $AdS_5 \times S^5$

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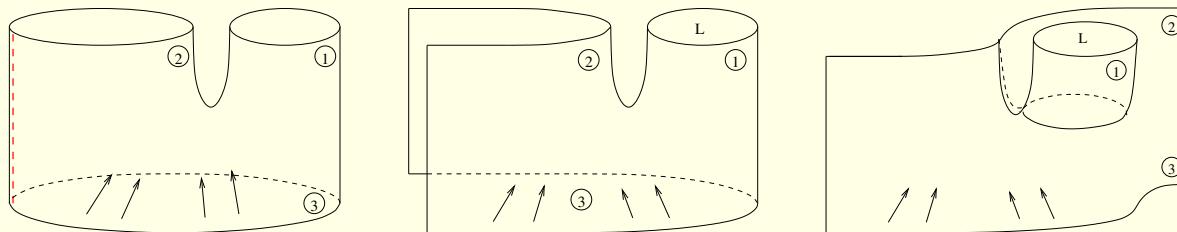


2 particles in string 3: $N_L(\theta_1, \theta_2) = N_L(\theta_2, \theta_1) = e^{-ip_1 L} N_L(\theta_2, \theta_1 - 2i\pi)$

kinematical singularity: $-i\text{Res}_\epsilon N_L(\theta + i\pi + \epsilon, \theta) = (1 - e^{ipL})$

Free massive boson: pp-wave limit of strings on $AdS_5 \times S^5$

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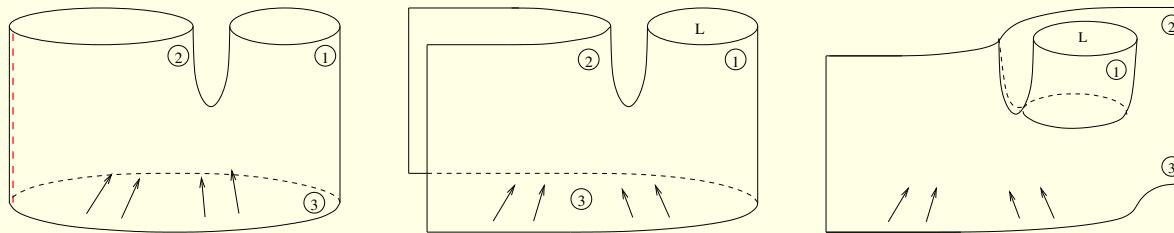
kinematical singularity: $-i\text{Res}_\epsilon N_L(\theta + i\pi + \epsilon, \theta) = (1 - e^{ipL})$

Solve the first two by:

$$N_L(\theta_1, \theta_2) = \frac{e^{-\theta_1 \frac{p_1}{2\pi} L} e^{-\theta_2 \frac{p_2}{2\pi} L}}{\cosh \frac{\theta_1 + \theta_2}{2}} n(\theta_1) n(\theta_2)$$

Free massive boson: pp-wave limit of strings on $AdS_5 \times S^5$

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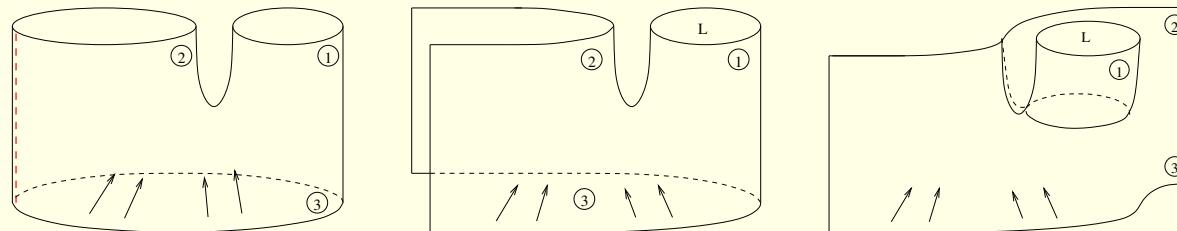
kin. singularity: $n(\theta) n(\theta + i\pi) = \sinh \theta \sin \frac{pL}{2}$ zeros at $\theta = \frac{2\pi n}{L}$ and $\theta = \frac{2\pi n}{L} + i\pi$

Free massive boson: pp-wave limit of strings on $AdS_5 \times S^5$

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$n(\theta) = \sinh \frac{\theta}{2} \sin \frac{pL}{2} \Gamma_{\frac{mL}{2\pi}}(m \sinh \theta)$ where Γ_μ removes zeros at $\theta = \frac{2\pi n}{L} + i\pi$

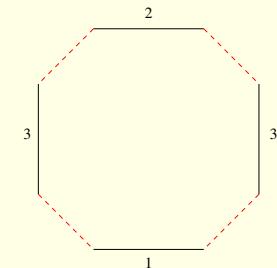
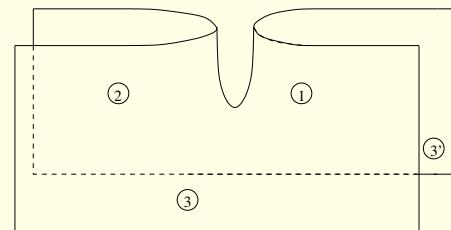
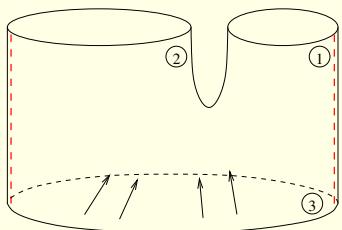
$$\Gamma_\mu(z) = z^{-1} e^{-\omega_z(\gamma + \log \frac{\mu}{2e})} \prod \frac{n}{\omega_n + \omega_z} e^{-\frac{\omega_n}{z}} \quad \text{and} \quad \omega_z = \sqrt{\mu^2 + z^2}$$

[Spradlin et al '02, Lucietti et al '03]

Free massive boson: summing up the octagons

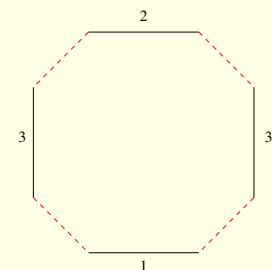
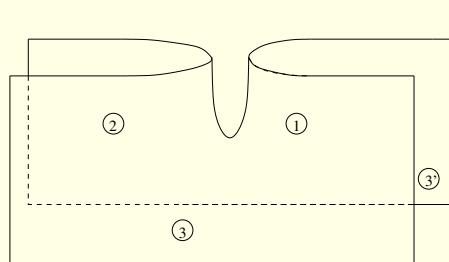
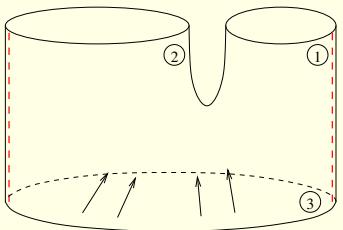
Free massive boson: summing up the octagons

2 cuts
octagon



Free massive boson: summing up the octagons

2 cuts
octagon



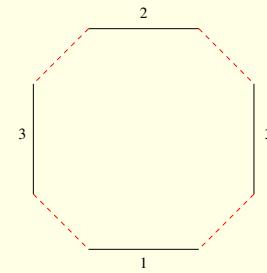
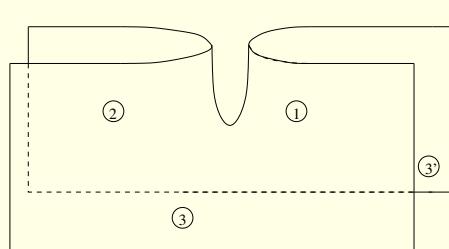
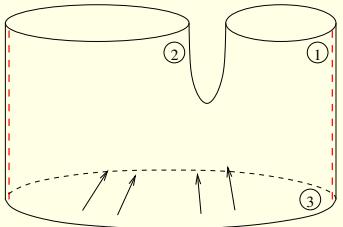
2 particles: $O(\theta_1, \theta_2) = O(\theta_2, \theta_1) = O(\theta_2, \theta_1 - 4i\pi)$

kinematical singularity: $-i\text{Res}_\epsilon O(\theta + i\pi + \epsilon, \theta) = 1$

$$\rightarrow O(\theta_1, \theta_2) = \frac{1}{\cosh \frac{\theta_1 - \theta_2}{2}}$$

Free massive boson: summing up the octagons

2 cuts
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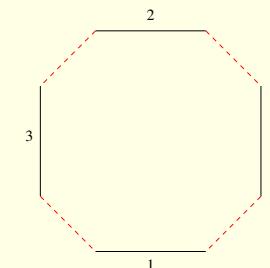
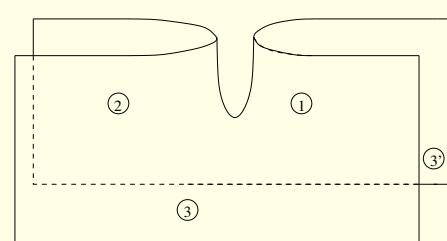
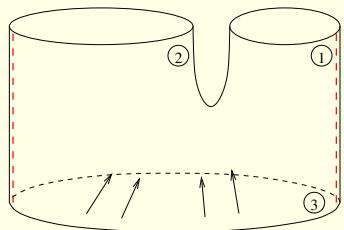
$$\rightarrow O(\theta_1, \theta_2) = \frac{1}{\cosh \frac{\theta_1 - \theta_2}{2}}$$

$$\text{Multiparticle solution: } O(1, 2, 3, 4) = O(1, 2)O(3, 4) + O(1, 3)O(2, 4) + O(1, 4)O(2, 3)$$

(Wick theorem)

Free massive boson: summing up the octagons

2 cuts
octagon

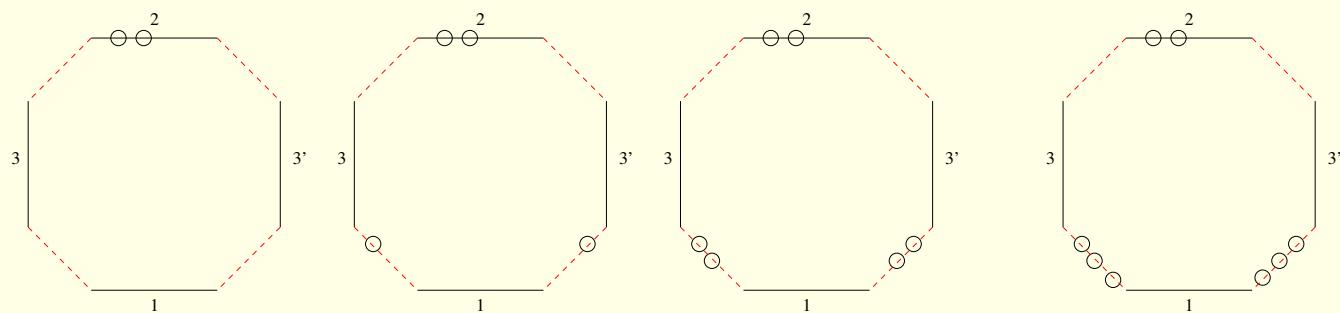


$$\text{2 particles: } O(\theta_1, \theta_2) = O(\theta_2, \theta_1) = O(\theta_2, \theta_1 - 4i\pi) \quad \rightarrow \quad O(\theta_1, \theta_2) = \frac{1}{\cosh \frac{\theta_1 - \theta_2}{2}}$$

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Summing up
virtual corrections



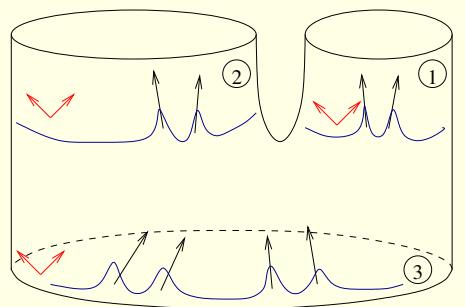
$$N_L(\theta_1, \theta_2) = O(\theta_1, \theta_2) + \int \frac{du}{2\pi} O(\theta_1, \theta_2, u - 3i\frac{\pi}{2}, u + 3i\frac{\pi}{2})_c e^{-m \cosh u L} + \dots$$

Agrees with the expansion of $N_L(\theta_1, \theta_2)!$

Conclusion

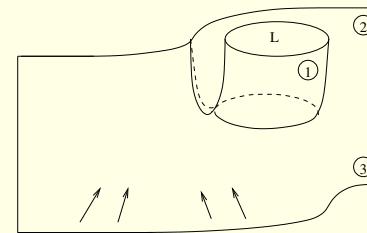
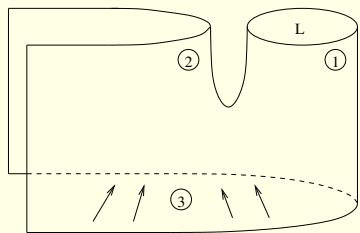
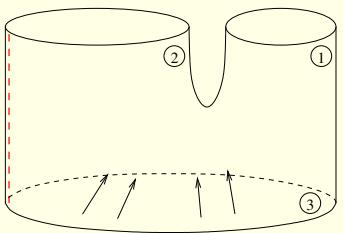
Conclusion

Ultimate goal:



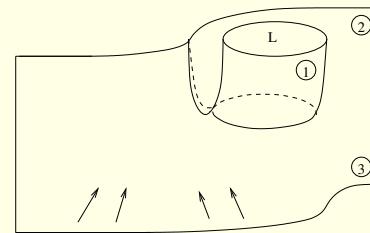
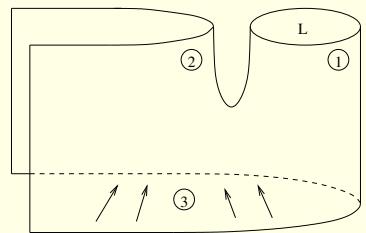
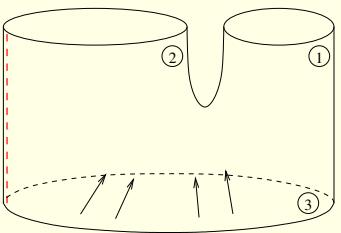
Conclusion

1 cut:
nonlocal
factors

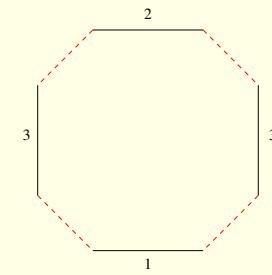
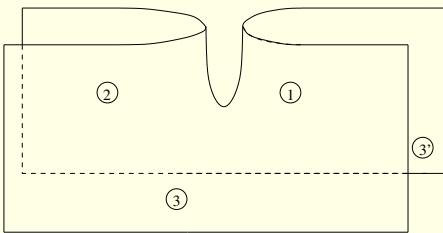
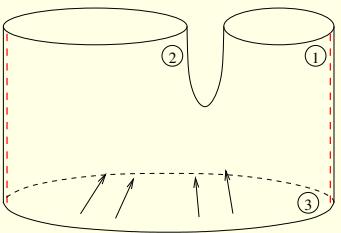


Conclusion

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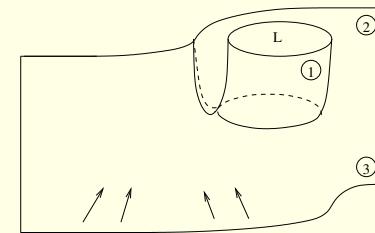
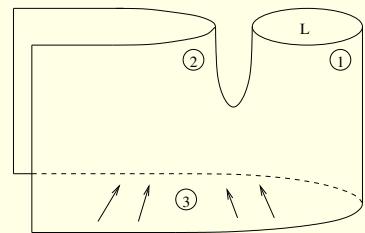
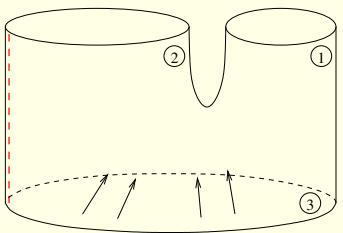


2 cuts
octagon

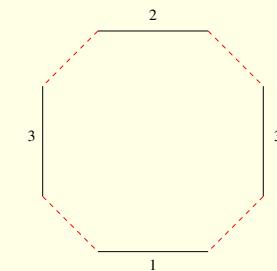
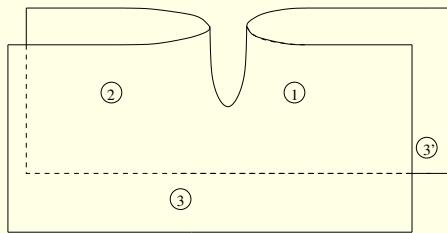
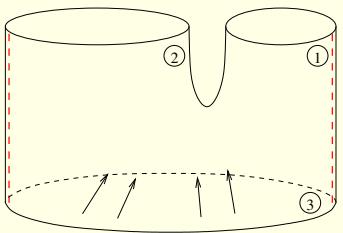


Conclusion

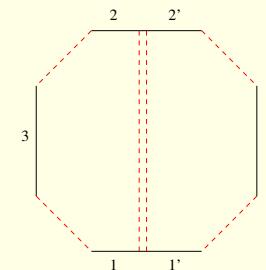
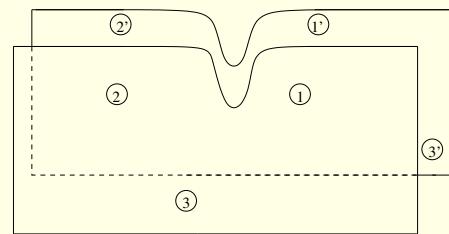
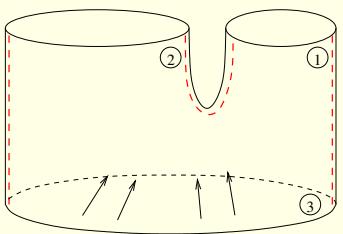
1 cut:
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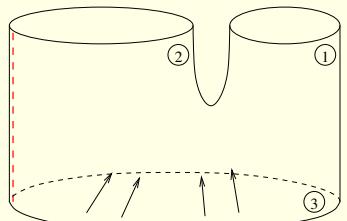


3 cuts
hexagon

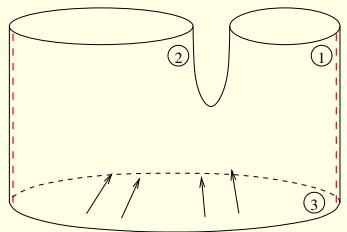


Conclusion

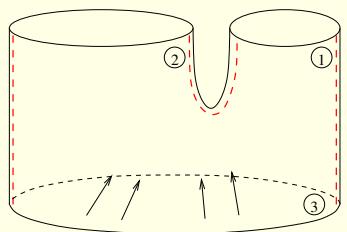
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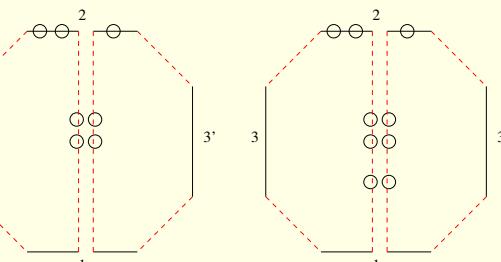
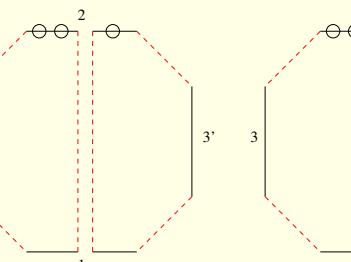
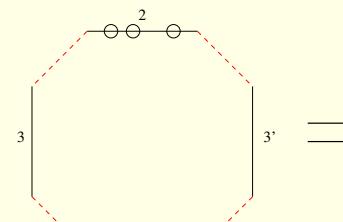
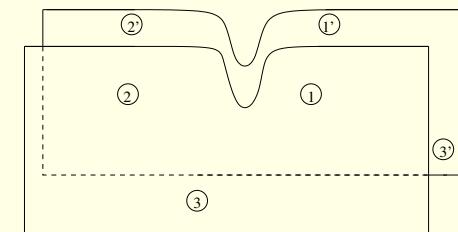
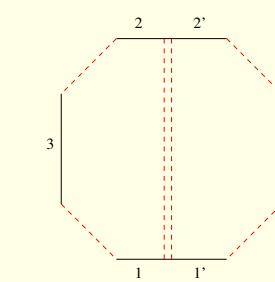
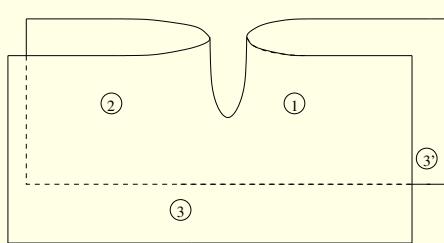
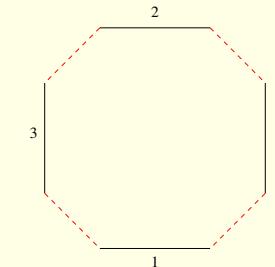
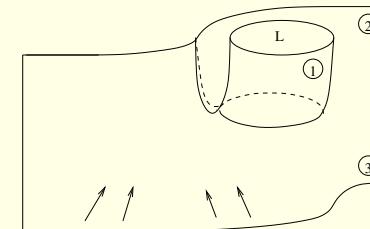
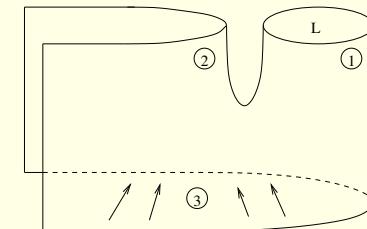
2 cuts
octagon



3 cuts
hexagon



sewing back

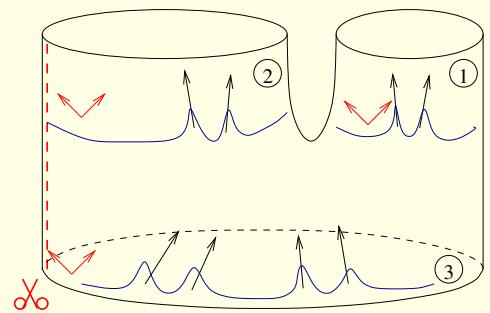


$$O(\theta_1, \theta_2, \theta_3) = h(\theta_1, \theta_2)h(\theta_3) + \dots + \int \frac{du}{2\pi} \mu(u) h(\theta_1, \theta_2, u - i\frac{\pi}{2})h(u + i\frac{\pi}{2}, \theta_3)e^{-E(u)l} +$$

The string vertex for $L_1 = 0$: diagonal form factor

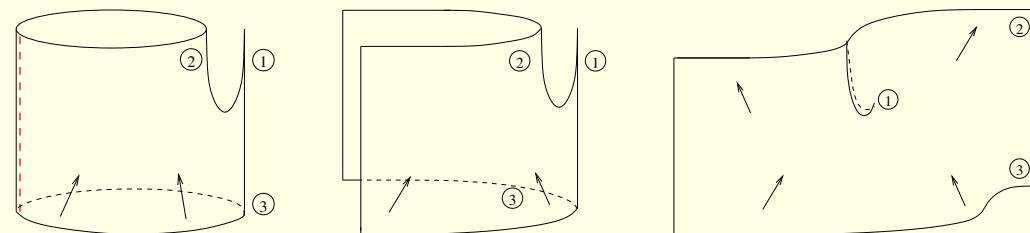
The string vertex for $L_1 = 0$: diagonal form factor

Decompatify string 2 & 3 but $L_1 = 0$:



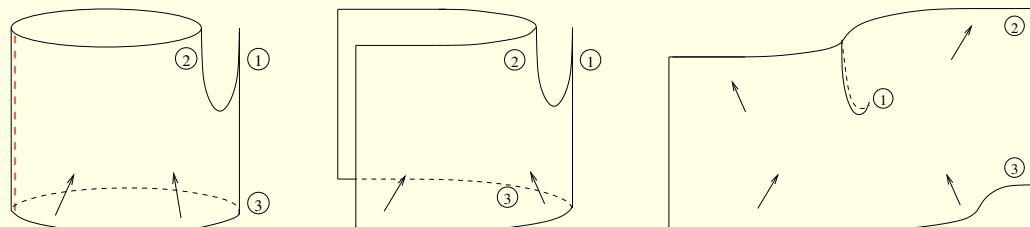
The string vertex for $L_1 = 0$: diagonal form factor

Decompectify string 2 & 3 and $L_1 = 0$:



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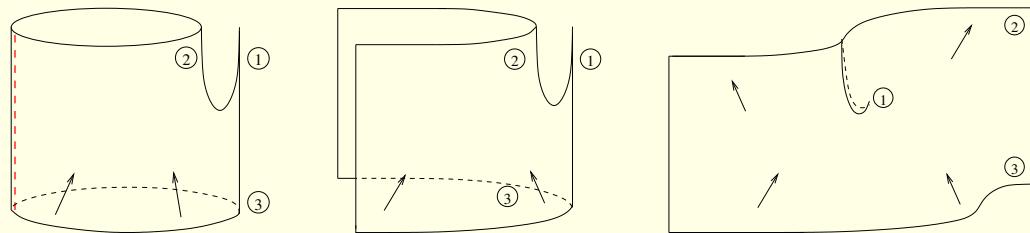
Local operator form factor equations:

$$N_0(\theta_1, \dots, \theta_n) = N_0(\theta_2, \dots, \theta_n, \theta_1 - 2i\pi) = S(\theta_i - \theta_{i+1}) N_0(\dots, \theta_{i+1}, \theta_i, \dots)$$

$$-i\text{Res}_{\theta'=\theta} N_0(\theta' + i\pi, \theta, \theta_1 \dots, \theta_n) = (1 - \prod_i S(\theta - \theta_i)) N_0(\theta_1, \dots, \theta_n)$$

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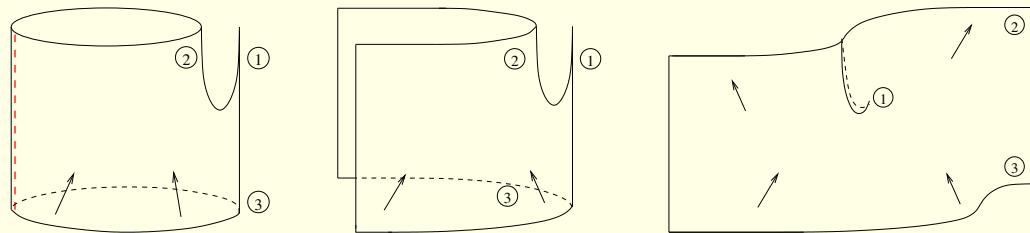
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HeavyHeavyLight 3pt function strong coupling prescription

[Costa et al., Zarembo]: $C_{HHL} = \int_{\text{world sheet}} \mathcal{V}(X[\text{heavy solution}(\sigma, \tau)]) d^2\sigma$

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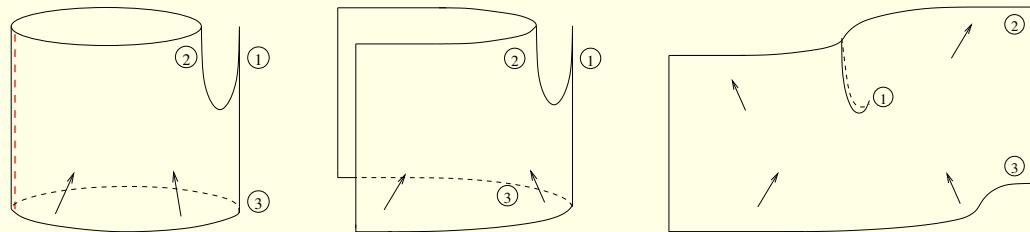
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for multiparticle state: $C_{HHL} = \int_{\text{moduli space } \{y_i\}} \mathcal{V}(X[\text{heavy solution}(y_i)]) d^n y$

The string vertex for $L_1 = 0$: diagonal form factor

Decompectify string 2 & 3 and $L_1 = 0$:



Local operator form factor equations:

$$N_0(\theta_1, \dots, \theta_n) = N_0(\theta_2, \dots, \theta_n, \theta_1 - 2i\pi) = S(\theta_i - \theta_{i+1}) N_0(\dots, \theta_{i+1}, \theta_i, \dots)$$

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HeavyHeavyLight 3pt function strong coupling prescription

[Costa et al., Zarembo]: $C_{HHL} = \int_{\text{world sheet}} \mathcal{V}(X[\text{heavy solution}(\sigma, \tau)]) d^2\sigma$

for multiparticle state: $C_{HHL} = \int_{\text{moduli space } \{y_i\}} \mathcal{V}(X[\text{heavy solution}(y_i)]) d^n y$

(classical) diagonal form factors: $\langle \theta_2, \theta_1 | \mathcal{V} | \theta_1, \theta_2 \rangle_L = \frac{F_2^s(\theta_1, \theta_2) + \rho_1(\theta_1)F_1^s(\theta_2) + \rho_1(\theta_2)F_1^s(\theta_1)}{\rho_2(\theta_1, \theta_2)}$

$e^{i\Phi_k} = 1$; $\Phi_k = p_k L - i \sum_{j:j \neq k} \log S(\theta_k, \theta_j)$; $\rho_n(\theta_1, \dots, \theta_n) = \det \left[\frac{\partial \Phi_j}{\partial \theta_i} \right]$

diagonal form factor $F_2^s(\theta_1, \theta_2) = \lim_{\epsilon \rightarrow 0} N_0(\bar{\theta}_2, \bar{\theta}_1, \theta_1 + \epsilon, \theta_2 + \epsilon)$

Explicitly checked at weak coupling [Hollo, Jiang, Petrovskii], AdS Form factors [McLoughlin, Klose]