On Quantum-Classical Hybrid Canonical Dynamics

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Abstract

Over the decades, desire of a hybrid of quantum and classical dynamics came from many fields spanning from quantum chemistry to cosmology, from foundations to open system theories, also from such special fields like quantum control. Koopman's quantum formalism of classical dynamics or Wigner's classical phase-space formalism of quantum dynamics are two opposite options to create the hybrid formalism. My topics is about a third option, kind of "in the middle". To construct natural coupling between a classical system's canonical formalism and a quantum system's operator formalism, Aleksandrov constructed a hybrid of the Poisson and Dirac brackets, Gerasimenko proposed an equivalent structure. This is remarkable and useful phenomenology but incorrect mathematically. Additional terms to the hybrid bracket can cure the defect, while the reversibility of the resulting hybrid dynamics becomes lost.

Desires of hybrid dynamics

| | CLASSICAL subSYSTEM | QUANTUM subSYSTEM |
|--------------|---------------------|--------------------|
| chemistry | nuclei | electrons |
| cosmology | gravity | matter |
| foundations | measuring device | measured system |
| open systems | reservoir | system of interest |
| control | measured signal | controlled system |

Three Main Options for Hybrid Dynamics

| CLASSICAL subSYSTEM | | QUANTUM subSYSTEM |
|------------------------|-------------------------------------|--------------------------|
| quantum-like formalism | $extend \! \Rightarrow \!$ | |
| Koopman | | |
| | ⇔extend | classical-like formalism |
| | | Wigner |
| classical formalism | \Rightarrow coupling \Leftarrow | quantum formalism |
| | Aleksandrov– | |
| | –Gerasimenko | |

Coupling Classical and Quantum Formalisms

| | C subSYSTEM | | Q subSYSTEM |
|-----------|---|---|--|
| State: | $\rho(q, p)$ | | $\hat{ ho}$ |
| Hamilton: | H(q, p) | | Ĥ |
| Motion: | $\dot{ ho} = \{H, \rho\}$ | | $\dot{\hat{ ho}}=-rac{i}{\hbar}[\hat{H},\hat{ ho}]$ |
| | Liouville eq. | | von Neumann eq. |
| | Poisson br. | | Dirac br. |
| | C formalism | \Rightarrow coupling \Leftarrow | Q formalism |
| State: | | $\widehat{ ho}(m{q},m{p})\equiv\widehat{ ho}$ | |
| Hamilton: | $\widehat{H}(q,p)\equiv \widehat{H}$ | | |
| Motion: | $\dot{\widehat{ ho}} = -rac{i}{\hbar} [\widehat{H}, \widehat{ ho}] + \operatorname{Herm} \{\widehat{H}, \widehat{ ho}\}$ | | |
| | AG hybrid eq. | | |
| | Aleksandrov br. | | |

AG Hybrid Dynamics 1981/82

$$\dot{\widehat{
ho}} = -rac{i}{\hbar} [\widehat{H}, \widehat{
ho}] + \operatorname{Herm} \{\widehat{H}, \widehat{
ho}\}$$

$$\dot{\hat{\rho}}(q,p) = - \frac{i}{\hbar} \left[\hat{H}(q,p), \hat{\rho}(q,p) \right] +$$

$$+ \operatorname{Herm} \left(\frac{\partial \hat{H}(q,p)}{\partial p} \frac{\partial \hat{\rho}(q,p)}{\partial q} - \frac{\partial \hat{H}(q,p)}{\partial q} \frac{\partial \hat{\rho}(q,p)}{\partial p} \right)$$

Useful effective dynamics. But inconsistent mathematically. D.-Gisin-Strunz 2000: 1D classical particle coupled to Pauli-spin:

$$\hat{H}(q,p) = H_{\it part}(q,p) + \hat{H}_{\it spin} + \kappa p \hat{\sigma}_{3}$$

AG hybrid eq. can destruct positivity $0 \leq \hat{\rho}(q, p)$. So what?

Quantum Dynamics with Two \hbar 's, D. 1995

- $\bullet\,$ Quantize the classical subsystem as well, but with $\hbar'\neq\hbar\,$
- Couple it to the quantum subsystem of interest. Educated Ansatz: generalization of Dirac bracket -(i/ħ)[.,.] for two ħ's.
- Nonunitary dynamics, like a quantum master eq. But!
- Incomplete Lindblad 1976 (GKLS 1976, in fact) master eq.
- Complete it! Add the minimum necessary new terms.
- Take $\hbar' \to 0$

Positivity Preserving Hybrid Dynamics, D. 1995

$$\begin{split} \hat{H}(q,p) &= \hat{H}_Q + H_C(q,p) + C(q,p)\hat{Q} \\ \hat{H} &= \hat{H}_Q + H_C + C\hat{Q} \end{split} \\ \dot{\hat{\rho}} &= -\frac{i}{\hbar}[\hat{H},\hat{\rho}] + \operatorname{Herm}\{\hat{H},\hat{\rho}\} - \frac{\lambda}{4\hbar^2}[\hat{Q},[\hat{Q},\hat{\rho}]] + \frac{1}{4\lambda}\{C,\{C,\hat{\rho}\}\} \\ & \operatorname{AG \ dynamics} \quad \operatorname{decoherence} \quad \operatorname{diffusion} \end{aligned} \\ \textbf{Least \ added \ noise:} \ (stregth \ of \ decoh.) \times (strength \ of \ diff.) = \operatorname{const.} \\ \textbf{Example: \ 1D \ classical \ particle \ coupled \ to \ Pauli-spin:} \end{split}$$

$$\hat{H}(q,p)=\hat{H}_Q+(p^2/2m)+\kappa p\hat{\sigma}_3$$

$$\dot{\hat{\rho}}(q,p) = -\frac{i}{\hbar} [\hat{H}_Q, \hat{\rho}(q,p)] + \frac{p}{m} \frac{\partial \hat{\rho}(q,p)}{\partial q} + \kappa \operatorname{Herm} \hat{\sigma}_3 \hat{\rho}(q,p) - \frac{\lambda \kappa^2}{4\hbar^2} [\hat{\sigma}_3, [\hat{\sigma}_3, \hat{\rho}(q,p)]] + \frac{\kappa^2}{4\lambda} \frac{\partial^2 \hat{\rho}(q,p)}{\partial q^2}$$

Positivity $0 \leq \hat{\rho}(q, p)$ guaranteed.

What's That, What's It Good or Not so Good for?

$$\widehat{H} = \widehat{H}_Q + H_C + C\,\widehat{Q}$$

$$\dot{\widehat{
ho}} = -rac{i}{\hbar}[\widehat{H},\widehat{
ho}] + \operatorname{Herm}\{\widehat{H},\widehat{
ho}\} - rac{\lambda}{4\hbar^2}[\hat{Q},[\hat{Q},\widehat{
ho}]] + rac{1}{4\lambda}\{C,\{C,\widehat{
ho}\}\}$$

- Special case of generic Hybrid Master Equations, generic unification of Pauli classical & GKLS quantum kinetic (master) equations, D. 2014
- Good, if interacion is written in the form $\sum C \hat{Q}$ (electrodynamics, weak gravity, many linearized couplings)
- Not so good if there is no distigushed expansion $\sum C_{\cdot}\hat{Q}_{\cdot}$ of coupling.

Three Main Options for Hybrid Dynamics — Comparisons?

| CLASSICAL subSYSTEM | | QUANTUM subSYSTEM |
|------------------------|-------------------------------------|--------------------------|
| quantum-like formalism | extend | |
| Koopman | | |
| | ⇔extend | classical-like formalism |
| | | Wigner |
| classical formalism | \Rightarrow coupling \Leftarrow | quantum formalism |
| | Aleksandrov– | |
| | –Gerasimenko | |
| | +decoherence | |
| | +diffusion | |
| | =consistency | |

Compare the these three! Take the simplest hybrid coupling:

 $\kappa p \times \sigma_3$

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