

A QUANTUM CORRECTION TERM <sup>①</sup>  
to  
FOKKER-PLANCK EQUATION

'e'1'3 e'1'd  
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Historic Context

FP to QuantumFP

Master Equation Form

Lindblad Theorem

Modified QFP

Motivation (heuristic)

Derivations (microscopic)

Summary

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# HISTORIC CONTEXT

(2)

Newton  $F_{\text{frict}} = -\eta v$

Einstein  $D_{pp} = \eta k_B T$   $\frac{\eta}{m} = \lambda$

	microscopic	phenomenologic	
classical	Boltzmann collision bath ⋮	Langevin s.d.e. ⋮	FP p.d.e. ⋮
quantum	Landau harmonic bath ⋮	Gardiner q.s.d.e. ⋮	Agarwal q.p.d.e. ⋮
	QFP: Redfield, Agarwal, Haake-Reibold, Caldeira-Leggett, Zurek, Hu-Paz-Zhao		

Kohen-Marston-Tannor (1997):

A Markovian theory of  $q$ -dissipation has been a much sought after goal of at least six communities: NMR,  $q$ -optics, condensed matter physicists, mathematical physicists, astrophysicists and condensed phase chemical physicists.

FP  $\rightarrow$  QFP p.d.e.

Classical dissipative motion

$W = W(q, p)$  phase-space density

$$\dot{W} = -\frac{p}{m} \frac{\partial W}{\partial q} + \frac{\partial V}{\partial q} \frac{\partial W}{\partial p} + \lambda \frac{\partial p W}{\partial p} +$$

$$+ D_{pp} \frac{\partial^2 W}{\partial p^2}$$

$\uparrow$  friction  
 $\uparrow$  p-diffusion

$\parallel$   
 $\lambda m k_B T$

Let  $W$  be Wigner-function of  $\hat{\rho}$

$$W(q, p) \Leftrightarrow \hat{\rho}$$

def:  $\hat{p}_c, \hat{q}_c$ :  $\hat{p}_c \hat{o} \equiv \frac{1}{2} \{\hat{p}, \hat{o}\}$ ,  $\hat{q}_c \hat{o} \equiv \frac{1}{2} \{\hat{q}, \hat{o}\}$

$$W(q, p) = \text{tr} \left( \delta(q - \hat{q}_c) \delta(p - \hat{p}_c) \hat{\rho} \right)$$

$$\hat{\rho} = W(\hat{q}_c, \hat{p}_c)$$

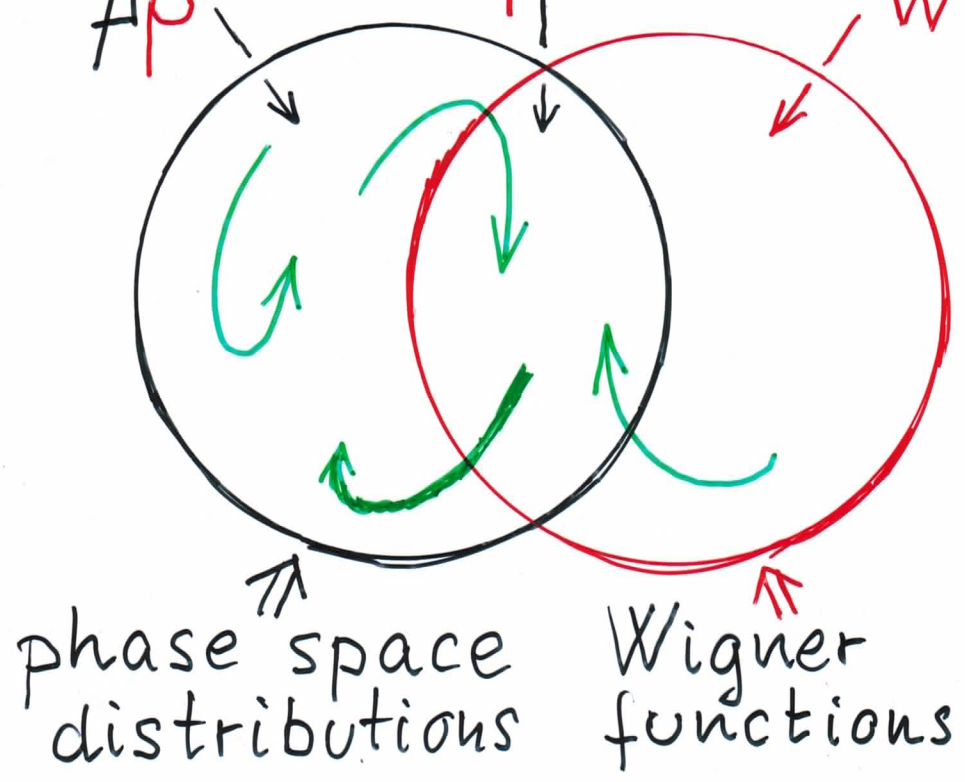
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note:  $\frac{\partial V}{\partial q} \frac{\partial W}{\partial p} \rightarrow \frac{2}{\hbar} V \left( q + \frac{i}{2} \hbar \frac{\partial}{\partial p} \right) W$

# THE BUG

FP p.d.e. solutions: 

$\nexists \hat{\rho}$   $\exists \hat{\rho}, W \geq 0$   $W \neq 0$





# QFP MASTER EQ

$$\hat{\rho} = W(\hat{q}_c, \hat{p}_c)$$

$$\dot{\hat{\rho}} = -\frac{i}{\hbar} \left[ \frac{\hat{p}^2}{2m} + V(\hat{q}), \hat{\rho} \right] - \frac{i}{2\hbar} \lambda [\hat{q}, \{\hat{p}, \hat{\rho}\}] - \frac{1}{\hbar^2} D_{PP} [\hat{q}, [\hat{q}, \hat{\rho}]]$$

Bug:  $\hat{\rho}_0 \xrightarrow{\text{green arrow}} \hat{\rho}_t \neq 0$

example:  $\hat{\rho}_0 = |\psi\rangle\langle\psi|$

$$\langle\psi|\hat{q}|\psi\rangle = 0, \quad \langle\psi|\hat{q}^2|\psi\rangle = \sigma^2$$

$$\text{let } |\varphi\rangle = \sigma^{-1} \hat{q} |\psi\rangle$$

$$\langle\varphi|\hat{\rho}_{\Delta t}|\varphi\rangle = \left( -\frac{\lambda}{2} + \frac{2}{\hbar^2} D_{PP} \sigma^2 \right) \Delta t$$

$\uparrow$   
 $\lambda m k_B T$

$\hat{\rho}_{\Delta t} \neq 0$  if

$$\sigma < \frac{\hbar}{\sqrt{2} \sqrt{2 m k_B T}}$$

# LINDBLAD, KOSSAKOWSKI-SUDARSHAN (6)

## -GORINI- THEOREM

special case:

$$\begin{aligned} \dot{\hat{\rho}} = & -\frac{i}{\hbar} [\hat{H}, \hat{\rho}] - \frac{i}{2\hbar} \lambda [\hat{q}, \{\hat{p}, \hat{\rho}\}] - \\ & - \frac{1}{\hbar^2} D_{pp} [\hat{q}, [\hat{q}, \hat{\rho}]] - \frac{1}{\hbar^2} D_{qq} [\hat{p}, [\hat{p}, \hat{\rho}]] - \\ & - \frac{2}{\hbar^2} D_{pq} [\hat{q}, [\hat{p}, \hat{\rho}]] \end{aligned}$$

is mathematically correct eq. iff:

$$\begin{pmatrix} D_{pp} & D_{qp} + \frac{i}{4}\lambda \\ D_{qp} - \frac{i}{4}\lambda & D_{qq} \end{pmatrix} \geq 0$$

In QFP:  $\begin{pmatrix} \lambda m k_B T & \frac{i}{4}\lambda \\ -\frac{i}{4}\lambda & 0 \end{pmatrix} \neq 0$

**Modify** QFP!

$$\uparrow \\ D_{qq}$$

$$D_{qq} = \kappa \frac{\lambda \hbar^2}{16 m k_B T}$$

$$\kappa \geq 1$$

# MOTIVATION (heuristic)

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Time coarse-graining of  
p-diffusion ( $D_{pp}$ ) yields  
q-diffusion ( $D_{qq}$ )

$$-\frac{\lambda m k_B T}{\hbar^2} [\hat{q}(\delta t), [\hat{q}(\delta t), \cdot]]$$

$$[\hat{q} + \frac{\hat{p}}{m} \delta t, [\hat{q} + \frac{\hat{p}}{m} \delta t, \cdot]]$$

$$[\hat{q}, [\hat{q}, \cdot]] + \frac{1}{m^2} (\delta t)^2 [\hat{p}, [\hat{p}, \cdot]]$$

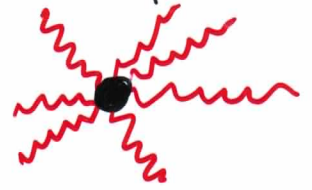
$$\delta t \sim \frac{\sqrt{\kappa}}{4} \frac{\hbar}{k_B T}$$

coarse-graining time-scale  $\sim$   
 $\sim \hbar / \text{thermal energy scale}$

# DERIVATION 1 (microscopic à la <sup>(8)</sup> Caldeira Leggett ...)

Coupling to **harmonic bath** of temp.  $T$

$$\hat{H}_{S+I} = \frac{\hat{P}^2}{2m} + V(\hat{q}) + \hat{q} \hat{Q}$$



$$\langle \hat{Q}(t) \rangle_B = 0, \quad \langle \hat{Q}(t) \hat{Q}(t') \rangle_B = \frac{1}{\hbar} \alpha(t-t')$$

$$\alpha_I(t) = -\frac{\lambda m}{\pi} \int_0^\Omega \omega \sin(\omega t) d\omega$$

$$\alpha_R(t) = \frac{\lambda m}{\pi} \int_0^\Omega \omega \coth\left(\frac{\hbar\omega}{2k_B T}\right) \cos(\omega t) d\omega$$

Markov approximation:

$$\alpha_I(t) \approx \lambda m \dot{\delta}(t) \Rightarrow -\frac{i}{2\hbar} \lambda [\hat{q}, \{\hat{P}, \hat{p}\}]$$

$$\alpha_R(t) \approx \frac{2\lambda m k_B T}{\hbar} \delta(t) - \frac{\lambda m \hbar}{6k_B T} \ddot{\delta}(t) + \dots$$



$$-\frac{1}{\hbar^2} \lambda m k_B T [\hat{q}, [\hat{q}, \hat{p}]]$$



$$-\frac{1}{\hbar^2} \frac{\lambda}{12m k_B T} [\hat{P}, [\hat{P}, \hat{p}]]$$

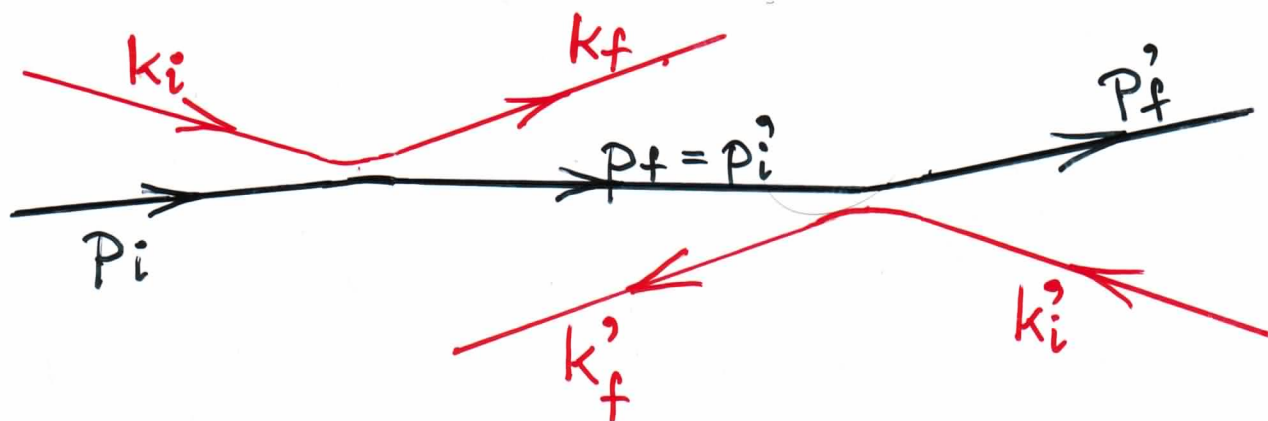
$$\left(\kappa = 4/3\right)$$

while  $D_{qp} \neq 0$



# DERIVATION 2 (microscopic <sup>⑨</sup> à la Boltzmann...)

Coupling to collision bath of temp.  $T$



Approximations:

$m \gg$  mass of bath molecules

small scattering angles  $\Theta$

time coarse-graining over many coll's

$$D_{pp} = \frac{1}{3} \underbrace{\langle \nu (k_f - k_i)^2 \rangle_{av}}_{\substack{\uparrow \\ \text{collision rate}}} = \left\{ \begin{array}{l} \text{squared-momentum} \\ \text{transfer rate} \end{array} \right.$$

$$\lambda = \frac{1}{m k_B T} D_{pp}$$

$$D_{qq} = \frac{1}{12} \left( \frac{\hbar}{m k_B T} \right)^2 \langle \nu k^2 \rangle_{av} \sim \left\{ \begin{array}{l} \text{scattered} \\ \text{energy flux} \end{array} \right.$$

$$\kappa = 4 \frac{\langle \nu k^2 \rangle_{av}}{\langle \nu (k_f - k_i)^2 \rangle_{av}} \sim \frac{4}{\langle \Theta^2 \rangle_{av}}$$

# SUMMARY

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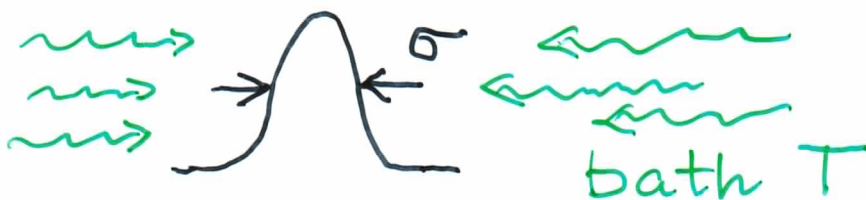
FP & naïve QFP:

$$\dot{W} = -\frac{p}{m} \frac{\partial W}{\partial q} + \frac{\partial V}{\partial q} \frac{\partial W}{\partial p} + \lambda \frac{\partial p W}{\partial p} + \lambda m k_B T \frac{\partial^2 W}{\partial p^2}$$

Modified QFP:

$$\dot{W} = \dots \text{naïve QFP} \dots + \kappa \hbar^2 \frac{\lambda}{16 m k_B T} \frac{\partial^2 W}{\partial q^2}$$

- Classical FP thermalizes always to Gibbs
- MQFP thermalizes free particle & harm. osc.
- MQFP  $\neq$  universal thermalizing engine
- MQFP is superior in numeric simulations
- Experimental significance?



$$\sigma < \frac{\hbar}{\sqrt{2 m k_B T}}$$