Temperature and heat conduction beyond Fourier law

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 - Guyer-Krumhansl
 - Jeffreys type
- 3. Experiments

Heat exchange experiment





http://remotelab.energia.bme.hu/index.php?page=thermocouple_remote_desc&lang=en



Model 1 (Newton, 3 parameters):

$$\dot{T} = -l(T - T_0)$$

	Estimate	Std. Error
1	0.02295	0.00014
T_0	52.52	0.02
T _{ini}	27.33	0.09



Model 1 (Newton, 3 parameters):

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Model 2 (Extended Newton, 5 parameters):

$$\tau \dot{T} + \dot{T} = -l(T - T_0)$$



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Model 2 (Extended Newton, 5 parameters):

$$\tau \vec{T} + \dot{T} = -l(T - T_0)$$

	Estimate	Std. Error
1	0.0026	0.0009
T_0	53.3	0.30
T _{ini}	26.98	0.05
τ	35.8	1.6
vT _{ini}	0.617	0.004



Oscillations in heat exchange:

parameters
 model - values
 micro - interpretation
 macro-meso mechanism

Fourier – local equilibrium (Eckart, 1940)

$$\rho \dot{e} + \partial^{i} q^{i} = 0$$

$$\rho \dot{s} + \partial^{i} J^{i} \ge 0$$

$$s(e), \quad J^{i} = \frac{1}{T} q^{j}$$

Entropy production:

$$\rho \dot{s} + \partial^i J^i = \rho \frac{ds}{de} \dot{e} + \partial^i \frac{q^i}{T} = -\frac{1}{T} \partial^i q^i + q^i \partial^i \frac{1}{T} + \frac{1}{T} \partial^i q^j = q^i \partial^i \frac{1}{T} \ge 0$$

Constitutive equations (isotropy):

$$q^{i} = L \partial^{i} \frac{1}{T} = -\frac{L}{T^{2}} \partial^{i} T = -\lambda \partial^{i} T, \qquad \lambda \ge 0$$
 Fourier law

<u>Cattaneo-Vernotte equation</u> (Gyarmati, 1977, modified)

$$\frac{\rho \dot{e} + \partial^{i} q^{i} = 0}{\rho \dot{s} + \partial^{i} J^{i} \ge 0} \qquad s \left(e - \frac{m}{2} q^{2} \right), \quad J^{i} = \frac{1}{T} q^{j}$$

Entropy production:

$$\rho \dot{s} + \partial^i J^i = -\frac{1}{T} \partial^i q^i - \frac{m}{T} q^i \dot{q}^i + \partial^i \frac{q^i}{T} = q^i \left(\partial^i \frac{1}{T} - \frac{m}{T} \dot{q}^i \right) \ge 0$$

Constitutive equations (isotropy):

$$q^{i} = L \left(\partial^{i} \frac{1}{T} - \frac{m}{T} \dot{q}^{i} \right) \Rightarrow \boxed{\frac{mL}{T} \dot{q}^{i} + q^{i}} = -\frac{L}{T^{2}} \partial^{i} T$$
Cattaneo-Vernotte

Heat conduction constitutive equations

$$\rho \dot{e} + \partial^i q^i = 0$$

$$q^{i} = -\lambda \partial^{i}T,$$

$$\overline{tq}^{i} + q^{i} = -\lambda \partial^{i}T,$$

$$\overline{tq}^{i} + q^{i} = -\lambda \partial^{i}T,$$

$$\overline{tq}^{i} + q^{i} = -\lambda \partial^{i}T + a_{1} \partial^{ij}q^{j} + a_{2} \partial^{jj}q^{i},$$

$$\overline{tq}^{i} + q^{i} = -\lambda \partial^{i}T + l\partial^{i}\dot{T},$$

$$\overline{tq}^{i} + q^{i} = -\lambda \partial^{i}T + l\partial^{i}\dot{T},$$

$$\overline{tq}^{i} = -\lambda \partial^{i}T + a_{2} \partial^{jj}q^{i}.$$

there are more...

Thermodynamic approach

vectorial internal variable and current multiplier (Nyíri 1990, Ván 2001)

$$\rho \dot{e} + \partial^{i} q^{i} = 0 \qquad s \left(e - \frac{m}{2} \xi^{2} \right), \quad J^{i} = B^{ij} q^{j}$$
$$\rho \dot{s} + \partial^{i} J^{i} \ge 0$$

Entropy production:

$$\rho \dot{s} + \partial^{i} J^{i} = -\frac{1}{T} \partial^{i} q^{i} - \frac{m\rho}{T} \xi^{i} \dot{\xi}^{i} + \partial^{i} \left(B^{ij} q^{j} \right) = \\ \partial^{i} q^{j} \left(\left(B^{ij} - \frac{1}{T} \delta^{ij} \right) + \left(\partial^{i} B^{ij} \right) q^{j} - \frac{m\rho}{T} \xi^{i} \dot{\xi}^{i} \ge 0$$

Constitutive equations (isotropy):

$$\begin{aligned} q^{i} &= l_{1}\partial^{j}B^{ji} - \hat{l}_{12}\xi^{i}, & \hat{l}_{12} = l_{12}\frac{\rho m}{T} \\ \dot{\xi}^{i} &= l_{21}\partial^{j}B^{ji} - \hat{l}_{2}\xi^{i}, & \hat{l}_{2} = l_{2}\frac{\rho m}{T} \\ B^{ij} - \frac{1}{T}\delta^{ij} &= k_{1}\partial^{i}q^{j} + k_{2}\partial^{j}q^{i} + k_{3}\partial^{k}q^{k}\delta^{ij}. \end{aligned} \qquad \begin{aligned} l_{1} \geq 0, & l_{2} \geq 0, \\ L &= l_{1}\hat{l}_{2} - l_{12}\hat{l}_{21} \geq 0 \\ k_{1} \geq 0, & k_{2} \geq 0, \\ k_{3} \geq 0 \end{aligned}$$

$$\tau \dot{q}^{i} + q^{i} = -\lambda_{1} \partial^{i} T - \lambda_{2} \partial^{i} \dot{T} + a_{1} \partial^{ij} q^{j} + a_{2} \partial^{jj} q^{i} + b_{1} \partial^{ij} \dot{q}^{j} + b_{2} \partial^{jj} \dot{q}^{i}$$

$$\tau = \frac{1}{l_2}, \quad \lambda_1 = \frac{L}{l_2 T^2}, \quad \lambda_2 = \frac{l_1}{l_2 T^2}, \quad \lambda_1 = \frac{L}{l_2 T^2},$$
$$a_1 = \frac{L}{l_2} (k_1 + k_3), \quad a_2 = \frac{L}{l_2} k_2, \quad b_1 = \frac{l_1}{l_2} (k_1 + k_3), \quad b_2 = \frac{l_1}{l_2} k_2,$$

<u>1+1 D:</u>

$$\rho c \dot{T} + q' = 0,$$

$$\tau \dot{q} + q = -\lambda_1 T' - \lambda_2 \dot{T}' + aq'' + b \dot{q}''.$$

$$\tau \ddot{T} + \dot{T} = \hat{\lambda}T'' + \hat{a}\dot{T}'' + a_1\ddot{T}''.$$

 $\dot{T} = \hat{\lambda}T''$ $\tau \ddot{T} + \dot{T} = \hat{\lambda}T''$ $\tau \ddot{T} + \dot{T} = \hat{\lambda}T'' + \hat{a}\dot{T}''$ $\tau \ddot{T} = \hat{\lambda}T'' + \hat{a}\dot{T}''$

Fourier

Cattaneo-Vernotte

Guyer and Krumhansl and Jeffreys type Green-Naghdi type

Calculations

'Meso' models

a) Jeffreys-type equation – heat separation

$$\hat{q}^{i} = -\hat{\lambda}\partial^{i}T \\ \tau \dot{\tilde{q}}^{i} + \tilde{q}^{i} = -\tilde{\lambda}\partial^{i}T$$

$$\hat{q}^{i} + \tilde{q}^{i} = q^{i}$$

1+1D:

$$\begin{aligned} \tau \dot{q} &= -\tilde{q} - \tilde{\lambda} T' - \tau \hat{\lambda} \dot{T}' = -q + \hat{q} - \tilde{\lambda} T' - \tau \hat{\lambda} \dot{T}' \\ \tau \dot{q} + q &= -(\tilde{\lambda} + \hat{\lambda}) T' - \tau \hat{\lambda} \dot{T}' \end{aligned} \qquad \text{Jeffreys} \end{aligned}$$

b) Jeffreys-type equation – dual phase lag

$$q^{i}(r,t+\tau_{1}) = -\lambda \partial^{i} T(r,t+\tau_{2})$$

Taylor series:

$$q^{i} + \tau_{1}\dot{q}^{i} = -\lambda\partial^{i}T - k\tau_{2}\partial^{i}\dot{T}$$
 Jeffreys

This is unacceptable.

c) Jeffreys-type equation – two steps

$$c_1 \dot{\mathbf{T}}_1 = -\partial^i q^i - g(T_1 - T_2),$$
$$q^i = -\lambda \partial^i T_1$$
$$c_2 \dot{\mathbf{T}}_2 = g(T_1 - T_2)$$

$$c_{1}\ddot{T}_{1} - \lambda \dot{T}_{1}'' + g\dot{T}_{1} = g\dot{T}_{2} = \frac{g^{2}}{c_{2}}(T_{1} - T_{2}) = \frac{g^{2}}{c_{2}}T_{1} - \frac{g}{c_{2}}(c_{1}\dot{T}_{1} - \lambda T_{1}'')$$

$$c_{1}\ddot{T}_{1} + g\left(1 + \frac{c_{1}}{c_{2}}\right)\dot{T}_{1} = \frac{g\lambda}{c_{2}}T_{1}'' + \lambda \dot{T}_{1}''$$
Jeffreys

Waves +... in heat conduction:

 thermodynamic frame nonlocal hierarchy (length scales)
 macro-meso mechanisms frame is satisfied

Heat conduction equations

	Memory	Nonlocality	Objectivity	Thermo- dynamics
Fourier	no	no	research	ok
Cattaneo-Vernotte	yes	no	research	ok
Guyer-Krumhansl Gyarmati-Nyíri, (linearized Boltzmann)	yes	yes	?	ok
Jeffreys type 1. internal variable, 2 .heat separation 3. dual phase lag, 4. two steps	yes	yes	?	ok
Ballistic-diffusive (Boltzmann split)	yes	yes	?	?

Experiments

Homogeneous inner structure – metals

- typical relaxation times: $\tau = 10^{-13}$ 10^{-17} s
- Cattaneo-Vernotte is accepted: ballistic phonons (nano- and microtechnology?)

Inhomogeneous inner structure

- typical relaxation times: $\tau = 10^{-3}$ 100 s
- experiments are not conclusive

Kaminski, 1990



Mitra-Kumar-Vedavarz-Moallemi, 1995



Herwig-Beckert, 2000

sand, different setupno effect

Roetzel-Putra-Das, 2003

 similar to Kaminski and Mitra et. al.
 small effect

Scott-Tilahun-Vick, 2009

repeating Kaminski and Mitra et. al.
no effect



Summary and conclusions

- Inertial and gradient effects in heat conduction
- Internal variables versus substructures (macro – micro)
 - black box universality
- No experiments for gradient effects (supressed waves?)

Ballistic-diffusive equation (Chen, 2001)

Separation of distribution function



Ballistic, analytic solution

$$, \quad f = \tilde{f} + \hat{f},$$

diffusive

1D:

$$\dot{u} + \partial_i q^i = \sigma_u, \quad u = \hat{u} + \widetilde{u}, \quad q^i = \widetilde{q}^i + \hat{q}^i$$
$$\tau \dot{\widetilde{u}} + \partial_i \widetilde{q}^i = -\widetilde{u},$$
$$\tau \dot{\widehat{q}}^i + \hat{q}^i = -\alpha \partial_i \hat{u}$$
?