

Can material time derivative be objective?

T. Matolcsi

*Dep. of Applied Analysis, Institute of Mathematics,
Eötvös Roland University, Budapest, Hungary*

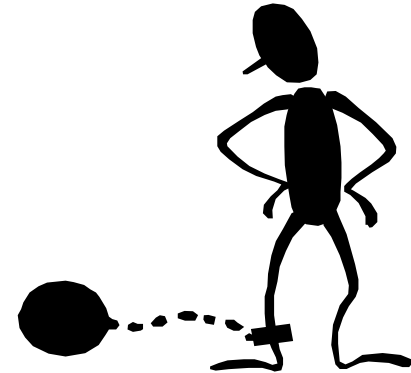
P. Ván

*Theoretical Department, Institute of Particle and Nuclear Physics,
Central Research Institute of Physics, Budapest, Hungary*

- Introduction
- About objectivity
- Covariant derivatives
- Material time derivative
- Jaumann derivative, etc...
- Conclusions



Classical
irreversible
thermodynamics



Local equilibrium (~ there is no microstructure)

Beyond local equilibrium:

Nonlocality in space
(structures)

Nonlocality in time
(memory and inertia)

constitutive space

Basic state space: $\mathbf{a} = (\dots)$

Nonlocality in space
(structures)

Nonlocality in time
(memory and inertia)

constitutive space
(weakly nonlocal)

$\nabla \mathbf{a}$
($\nabla^2 \mathbf{a}, \dots, \nabla^n \mathbf{a}$)

???

$\dot{\mathbf{a}}$
($\dot{\mathbf{a}}, \dots, \mathbf{a}^{(n\cdot)}$)

Nonlocality in spacetime

Rheology

Jaumann (1911)

Oldroyd (1949, ...)

...

Thermodynamic theory:

Kluitenberg (1962, ...),

Kluitenberg, Ciancio and Restuccia (1978, ...)

Verhás (1977, ..., 1998)

Thermodynamic theory with co-rotational time derivatives.

Experimental proof and prediction:

- viscometric functions of shear hysteresis
- instability of the flow !!

Material frame indifference

Noll (1958), Truesdell and Noll (1965)

Müller (1972, ...) ([kinetic theory](#))

Edelen and McLennan (1973)

Bampi and Morro (1980)

Ryskin (1985, ...)

Lebon and Boukary (1988)

Massoudi (2002) ([multiphase flow](#))

Speziale (1981, ..., 1998), ([turbulence](#))

Murdoch (1983, ..., 2005) and Liu (2005)

Muschik (1977, ..., 1998), Muschik and Restuccia (2002)

.....

Objectivity

About objectivity:

$$\hat{t} = t, \quad \hat{\mathbf{x}} = \mathbf{h}(t) + \mathbf{Q}(t)\mathbf{x} \quad \text{Noll (1958)}$$

$$\hat{x}^0 = x^0, \quad \hat{x}^\alpha = h^\alpha + Q^\alpha_\beta x^\beta, \quad \alpha, \beta = 1, 2, 3$$

$$\hat{x}^i = \hat{x}^i(x), \quad i = 0, 1, 2, 3$$

$c^i = (c^0, \mathbf{c})$ is a four dimensional objective vector, if

$$\hat{c}^i = \hat{J}^i_j c^j, \quad \text{where} \quad \hat{J}^i_j = \frac{\partial \hat{x}^i}{\partial x^j}$$

$$\hat{J} = \begin{pmatrix} 1 & \mathbf{0} \\ \dot{\mathbf{h}} + \dot{\mathbf{Q}}\mathbf{x} & \mathbf{Q} \end{pmatrix}$$

$$\begin{pmatrix} \hat{c}^0 \\ \hat{\mathbf{c}} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ \dot{\mathbf{h}} + \dot{\mathbf{Q}}\mathbf{x} & \mathbf{Q} \end{pmatrix} \begin{pmatrix} c^0 \\ \mathbf{c} \end{pmatrix} = \begin{pmatrix} c^0 \\ (\dot{\mathbf{h}} + \dot{\mathbf{Q}}\mathbf{x})c^0 + \mathbf{Q}\mathbf{c} \end{pmatrix}$$

Spec. 1: $c^0 = 0 \Rightarrow \hat{\mathbf{c}} = \mathbf{Q}\mathbf{c}$

Spec. 2: motion $\mathbf{r}(t) \Rightarrow \dot{\mathbf{r}} = \mathbf{v}$

$$\hat{\mathbf{r}} = \mathbf{h} + \mathbf{Q}\mathbf{r} \Rightarrow \dot{\hat{\mathbf{r}}} = \hat{\mathbf{v}} = \dot{\mathbf{h}} + \dot{\mathbf{Q}}\mathbf{r} + \mathbf{Q}\dot{\mathbf{r}}$$

$(t, \mathbf{r}(t)) \Rightarrow (1, \mathbf{v}(t))$ is an objective four vector

$$\begin{pmatrix} 1 \\ \hat{\mathbf{v}} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ \dot{\mathbf{h}} + \dot{\mathbf{Q}}\mathbf{x} & \mathbf{Q} \end{pmatrix} \begin{pmatrix} 1 \\ \mathbf{v} \end{pmatrix}$$

Covariant derivatives:

as the spacetime is flat there is a distinguished one.

$$Da \sim D_i a = \partial_i a \quad \text{covector field}$$

$$Dc \sim D_j c^i = \partial_j c^i + \Gamma_{jk}^i c^k \quad \text{mixed tensor field}$$

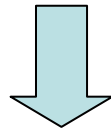
The coordinates of the covariant derivative of a vector field do not equal the partial derivatives of the vector field if the coordinatization is not linear.

If \hat{x} are inertial coordinates, the Christoffel symbol with respect to the coordinates x has the form:

$$\Gamma_{jk}^i = \frac{\partial^2 \hat{x}^m}{\partial x^j \partial x^k} \frac{\partial x^i}{\partial \hat{x}^m} = - \frac{\partial^2 x^i}{\partial \hat{x}^m \partial \hat{x}^l} \frac{\partial \hat{x}^m}{\partial x^j} \frac{\partial \hat{x}^l}{\partial x^k}$$

$$\hat{t} = t, \quad \hat{\mathbf{x}} = \mathbf{h}(t) + \mathbf{Q}(t)\mathbf{x}$$

$$\hat{x}^0 = x^0, \quad \hat{x}^\alpha = h^\alpha + Q^\alpha_\beta x^\beta, \quad \alpha, \beta = 1, 2, 3$$



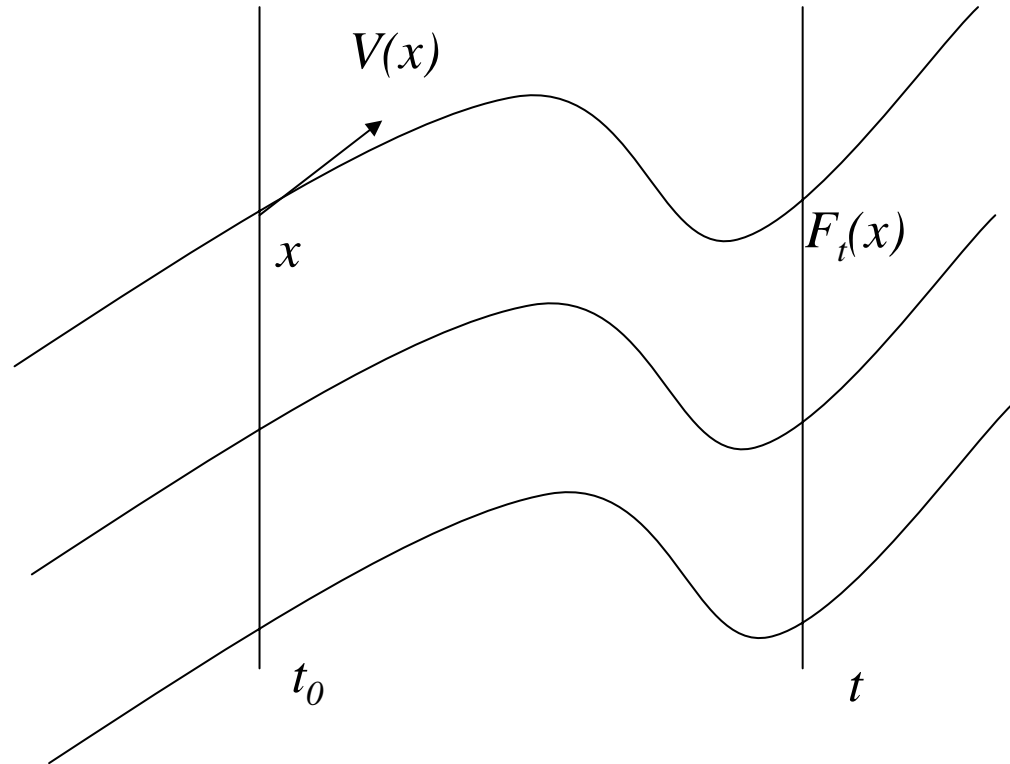
$$\Gamma_{jk}^0 = 0, \quad \Gamma_{0\beta}^\alpha = \Gamma_{\beta 0}^\alpha = \Omega^\alpha_\beta, \quad \Gamma_{\beta\gamma}^\alpha = 0,$$

where $\boldsymbol{\Omega} = \mathbf{Q}^{-1}\dot{\mathbf{Q}}$ is the angular velocity of the observer

$$\Gamma_{00}^\alpha = (\mathbf{Q}^{-1}(\ddot{\mathbf{h}} + \ddot{\mathbf{Q}}\mathbf{x}))^\alpha = (\mathbf{Q}^{-1}\ddot{\mathbf{h}} + (\dot{\boldsymbol{\Omega}} + \boldsymbol{\Omega}\boldsymbol{\Omega})\mathbf{x})^\alpha$$

Material time derivative:

Flow generated by
a vector field V .

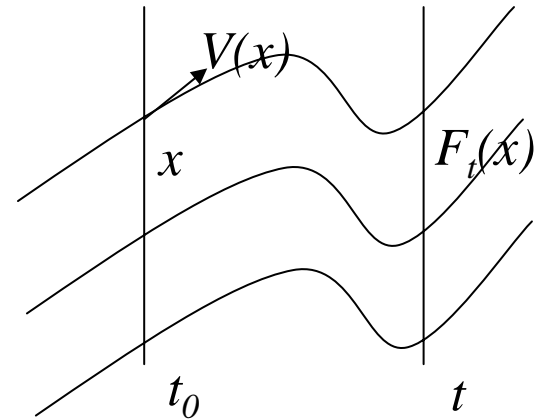


$F_t(x)$ is the point at time t of the integral curve V passing through x .

$\Phi(F_t(x))$ is the change of Φ along the integral curve.

$$\frac{d}{dt} \Phi(F_t(x)) = D_V \Phi(F_t(x))$$

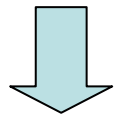
$D_V \Phi = V \cdot D\Phi$ is the covariant derivative of Φ according to V .



$$V = (1, \mathbf{v}) !$$

Spec. 1: Φ is a scalar

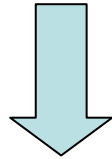
$$D_V a = V^j D_j a = V^j \partial_j a$$



$$D_V a = (\partial_0 + \mathbf{v} \cdot \nabla) a = \dot{a} \quad \text{substantial time derivative}$$

Spec. 2: $\Phi = c = (0, \mathbf{c})$ is a spacelike vector field

$$(D_V c)^i = V^j D_j c^i = V^j (\partial_j c^i + \Gamma_{jk}^i c^k)$$



$$D_V \mathbf{c} = (\partial_0 + \mathbf{v} \cdot \nabla + \mathbf{\Omega}) \mathbf{c}$$

The material time derivative of a vector – even if it is spacelike – is not given by the substantial time derivative.

Jaumann, upper convected, etc... derivatives:

In our formalism: ad-hoc rules to eliminate the Christoffel symbols.

For example:

$$V^j D_j c^i = V^j (\partial_j c^i + \Gamma_{jk}^i c^k)$$

$$c^j D_j V^i = c^j (\partial_j V^i + \Gamma_{jk}^i V^k)$$

$$\Downarrow \Gamma_{jk}^i = \Gamma_{kj}^i$$

$$V^j D_j c^i - c^j D_j V^i = V^j \partial_j c^i - c^j \partial_j V^i$$

for a spacelike vector $(0, \mathbf{c})$

$$V^j \partial_j c^i - c^j \partial_j V^i = (\partial_0 + \mathbf{v} \cdot \nabla) \mathbf{c} - \mathbf{c} \cdot \nabla \mathbf{v}$$

upper convected (contravariant) time derivative

One can get similarly Jaumann, lower convected, etc...

Conclusions:

- Objectivity has to be extended to a four dimensional setting.
- Four dimensional covariant differentiation is fundamental in non-relativistic spacetime. The essential part of the Christoffel symbol is the angular velocity of the observer.
- Partial derivatives are not objective. A number of problems arise from this fact.
- Material time derivative can be defined uniquely. Its expression is different for fields of different tensorial order.

space + time \neq spacetime

Thank you for your
attention.

